

## THE VENUSIAN ENVIRONMENT

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## **1 INTRODUCTION**

This report presents the first stage of results of a study of the atmospheric environment on Venus. The purpose of the study is to combine existing knowledge derived from measurements, theoretical studies and numerical general circulation models into a single model which can be used to specify conditions such as temperature, pressure, wind, humidity, and dust loading over a wide altitude range at any location, time of year, and time of day on Venus. This environmental model is intended for use in the design of future missions to Venus. The European Space Agency has already studied three such projects, 'Venus Express' orbiter, 'Lavoisier' entry probe, and 'Venus Sample Return'. The actual implementation of these or something similar is an increasingly likely prospect in the next decade.

The physical and dynamical properties of the atmosphere clearly are a major factor to be taken into account in designing hardware which can enter the atmosphere of Venus and perform properly under the extreme conditions which are known to prevail there. Other aspects to be studied with the aid of the model include the accuracy with which a particular design of vehicle can optimise its orbit using aerobraking, again taking into account factors such as the large day-night density variations in Venus' upper atmosphere. The model therefore has to include not only the most probable description of the key atmospheric parameters, but also reliable estimates of the uncertainty in, and natural variance of, these quantities.

The following sections contain a review of the existing knowledge of Venus and its atmosphere, and of how this knowledge was acquired. Section 2 provides, by way of introduction, some discussion of the main factors which design engineers need to take into account in planning a mission to enter Venus orbit safely and to operate inside the atmosphere and on the surface. The gross properties of the Venusian system are reviewed in Section 3, making appropriate comparisons with the other terrestrial planets, while sections 4-6 review in more detail the surface environment, atmospheric properties and global circulation respectively. Some background information on model formulation and presentation of outputs is also presented in Section 7 as a basis for the next stage of the study, the production of a new and improved model. The type of database which could result, and its applications, are covered in section 8 and the prospects for a successful European Venus modelling project are summarised in section 9. Finally, a comprehensive set of references to previous work is presented in section 10.

## **2 SPACECRAFT ENVIRONMENTAL FACTORS**

This section discusses environmental factors which may be of importance in planning a spacecraft mission which involves entering the atmosphere of Venus. The factors which affect the design of an entry probe fall into the following main categories: deployment variables, such as entry angle and parachute size, etc.; factors which affect the descent and landing site, such as pressure and wind speed; and factors which affect survival and the implementation of scientific experiments on the surface, such as temperature and pressure. The factors of importance for aerobraking into Venus orbit are principally the density, temperature and winds above 100 km altitude. These are considered in turn in the following sections.

### ***2.1 ENTRY PROBES and LANDERS***

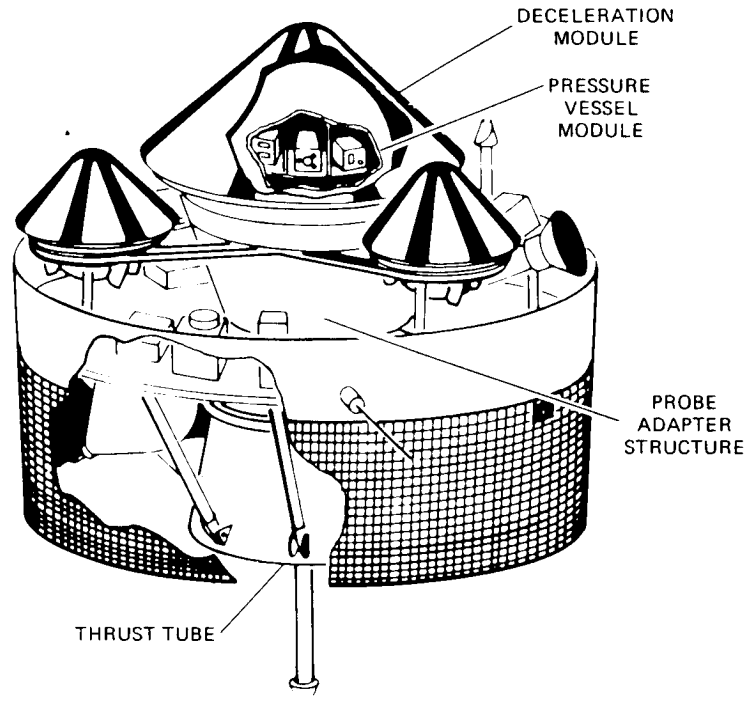


Figure 1. The Pioneer Venus Multiprobe spacecraft, showing the positioning of the cone-shaped probes on the platform of the multiprobe bus (NASA).

### Entry And Descent Phase

To perform a detailed entry and descent trajectory analysis will require vertical profiles of atmospheric temperature, pressure and wind as a function of height at all geographical locations under consideration for landing. During the entry, an entry probe experiences hypersonic, transonic and subsonic flight, starting out at a velocity likely to be in the range 5 to 15 km/s (the velocity was 11.6 km/s for the Pioneer Venus probes). A key parameter in designing the entry trajectory is determining the sonic altitude, or height above the surface where the entry vehicle decelerates through the sound barrier. This is likely to be in the range from 70 to 100 km.

The primary factors controlling behaviour on entry for a given set of conditions are the shape of the module, factors such as its mass distribution, and the design of the parachutes. The former is likely to be a basic sphere-cone configuration with a cone angle in the region of 45 to 60 degrees, for maximum drag and high stability against tilting or tumbling during descent. The detailed design of the entry module must also take into account factors such as the heating effects of atmospheric drag, which affects the amount of insulation and ablative material which must be carried, and the constraints imposed by the accommodation available for the landers on the launch and interplanetary cruise vehicles.

The flow field around the module must be known in order to compute the heating rates and the pressure distributions, which are needed to design the heat shield and the structure. In studies carried out for the NASA Pioneer project, entry was assumed to commence at a height of 200 km. The initial deceleration was accomplished using an aeroshell with an ablative heat shield of carbon phenolic bonded to its front face. A peak deceleration of 280g occurred soon after entry. At an altitude, which was not measured but can be estimated to be approximately 70km, the parachute was deployed. A pilot chute was mortar-fired from a compartment in the aeroshell at the same time as a pyrotechnic device was fired to release the main parachute cover. The pilot chute then pulled the main chute from its compartment; as soon as

deployment was complete and stability obtained the aeroshell was released and allowed to fall away.

Testing of the Pioneer entry system in the mid-1970s, was mainly empirical, using wind tunnels and drop-tests from aircraft. The most recent missions to enter Venus' atmosphere were VEGA 1 and 2 in 1985. Each of these identical spacecraft included a balloon station for atmospheric measurements, as well as a lander similar in design to earlier Venera missions. The lander separated from the Vega spacecraft two days before arrival at Venus and entered the planet's atmosphere on a ballistic path. (The flyby spacecraft was then retargetted using Venus gravity assistance to intercept Comet Halley in March 1986.) Vega 2 entered the atmosphere on 15 June 1985 at 01:59:30 UT at 10.80 km/s with an entry angle of 19.08 degrees. The pilot parachute was deployed at 02:00:05 UT at an altitude of 65 km and the braking parachute opened 11 seconds later at 64.5 km. The upper heat protection hemisphere was released at this time and the lower hemisphere 4 seconds later at 64.2 km. The upper hemisphere contained the deployment system for the balloon aerostat, which was deployed at an altitude of 54 km. The parachute was released at 02:09:15 at 47 km. After this the lander was allowed to aerobrake through the thick Venus atmosphere, with drag devices minimizing vibrations and spin and providing stability. The lander touched down at 03:00:50 UT on 15 June 1985 at 8.5 S, 164.5 E, in eastern Aphrodite Terra.

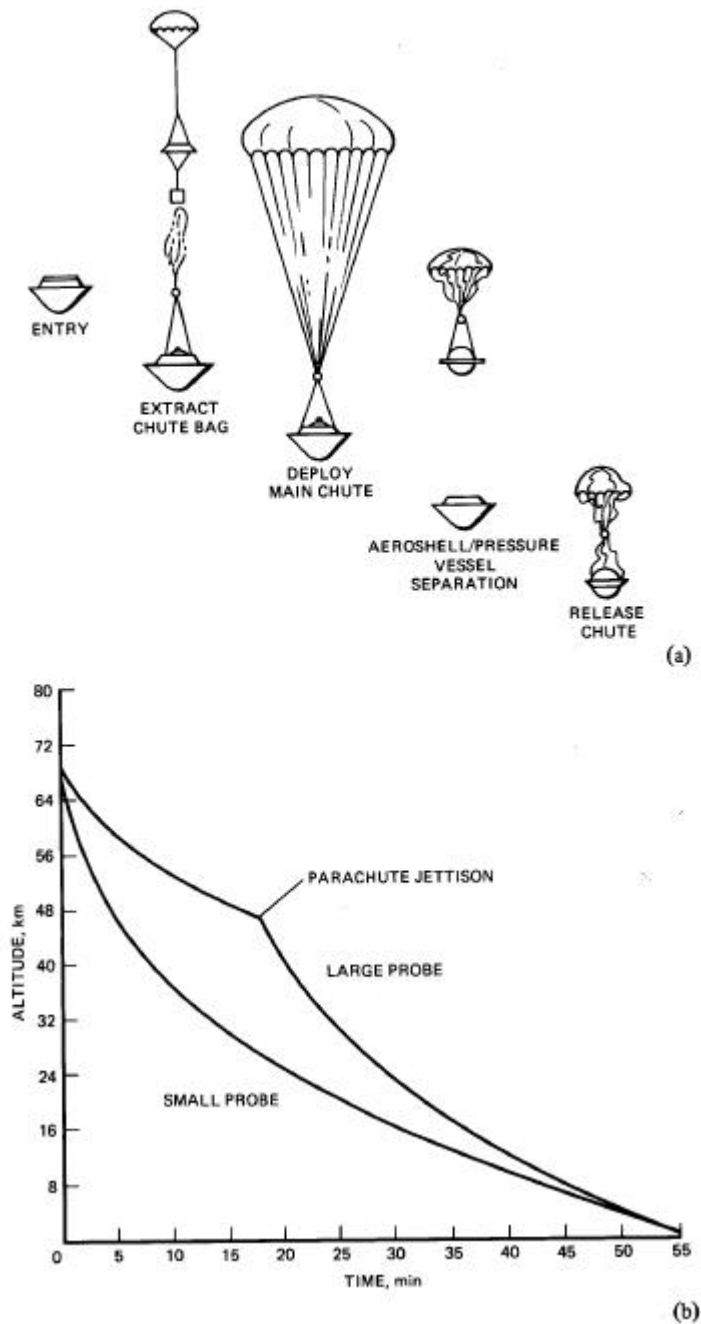


Fig.2. (a) The sequence of events in the descent of the Pioneer Venus probes. (b) Estimated height versus time for the Pioneer large and small probes; both took the same total time to reach the surface.

After release from the heat protection hemisphere, the balloon deployed a two-stage parachute, and then unfolded and inflated. The 3.4 meter diameter balloon supported a total mass of 25-kg. A 5-kg payload hung suspended 12 meters below the balloon. It floated at approximately 50 km altitude in the middle, most active layer of the Venus three-tiered cloud system. Data from the balloon instruments were transmitted directly to Earth for the 47-hr lifetime of the mission. (The batteries had a lifetime of 60 hrs.) Very long baseline interferometry was used to track the motion of the balloon to provide the wind velocity in the clouds. The tracking was to be done by a 6-station network on Soviet territory and by a



network of 12 stations distributed world-wide (organized by France and the NASA Deep Space Network). After two days the balloon entered the dayside of Venus and expanded and burst due to solar heating.

The most recent design of an entry system for Venus was that undertaken by ESA in 1988 as part of a feasibility study for a Venus Sample Return Mission. This adopted the basic Pioneer Venus aerobraking design, with a 45 deg conical heat shield 4m in diameter and an entry mass of 4000kg, and a parachute system adapted from Huygens. A light cross-canopy drogue chute 1.7 m in diameter is deployed at Mach 0.8 to extract the back cover and the main chute, the latter 4.6m in diameter. A total entry time of 4000 s was calculated, of which just under 1 hr is spent below a height of 50 km.

### **Landing And Surface Operation**

The Pioneer Venus probes were not designed for a soft landing on Venus, although one of them did continue to transmit for over an hour following impact. The Venera and VEGA landers used a toroidal system designed to absorb shock on landing. The most recent (sample return) study envisaged the deployment of a balloon (later to be used for the first stage of the ascent) during the final stages of the descent, to reduce the impact velocity to approximately 8 m/s. The lower part of the lander structure, exposed after release of the heatshield, is designed to be collapsible in order to absorb any remaining landing shock. It was calculated that 10 cm of crushable honeycomb material would be sufficient for the low impact velocity anticipated.

Once landed, the vehicle will need to survive the ambient conditions, in particular the very high temperature and pressure environment (720 K and 92 bar respectively). Depending on the operations to be performed at the surface, providing a suitable pressure vessel is not expected to be a major problem, although the mass implications are important. Winds are thought to be slight, at about 1 m/s typically. Temperature is the most serious problem; in particular the effect on electronic subsystems inside the spacecraft. Although high-temperature electronics have been under development for some time, it is still likely that any mission to the surface of Venus in the foreseeable future will be restricted to the use of semiconductors with a maximum tolerance of about 40 °C. Any kind of active refrigeration on the surface of Venus is similarly impractical at present and the only reasonable solution is temporary protection from the environment using insulating and heat-absorbing materials.

The ESA study envisioned a cocoon of high-temperature multi-insulation, manufactured by stacking and sewing together crinkled reflective foils separated by ceramic fabric. In the more external part metallic, ceramic or PBO (Poly-Benzobis-Oxazole) films can be used directly, kapton, aramid or other polymers need to be combined with, for example, silica fabrics. A computation based on a selected blanket 47 Kg/m<sup>3</sup> density gives a maximum internal temperature of 65 °C after 5 hours exposition to an environment with constant temperature of 450 °C.

The other element of the thermal protection system is the cooling system of the external gas injected into the cocoon. This is designed to take advantage of the presence of the helium, needed to inflate the balloon to cool down the incoming gas through a heat exchanger set in correspondence of the cocoon opening.

As the internal space is filled by the external CO<sub>2</sub>, during the descent time, the mass flow rate coming inside is low and so is the heat power to be removed. This permits utilisation of a compact and relatively lightweight heat exchanger with the helium flowing between the tanks and the balloon and the CO<sub>2</sub> being injected in the cocoon after cooling.

Electronics survival can be further enhanced by increasing the thermal inertia by the use of Phase Change Materials (PCM) in order to absorb the additional heat dissipated when the

components are in operation. Materials are available that have high transformation temperatures, high latent heat and low density, combined with low volumetric change, although the last criterion limits the choice to solid to liquid and solid to solid (chemical transitions) transformations.

The material with the best performance is phosphonium chloride but very few data are available on it; on the other hand, Lithium Nitrate Trihydrate (density 1550 Kg/m<sup>3</sup>) is well known and experimental data is available in the open literature. Disregarding the thermal ballast effect of the PCM itself, to keep within its temperature limits each component during the descent/surface/ascent phases, 0.036 Kg of material for each watt dissipated by the electronics are needed for 5 hours survival time. The total mass is then 9 Kg. Many configurations are possible. The simplest one is to directly attach a PCM metallic container to each electronic box by means of high conductance devices. Another possibility is to use a honeycomb support plate filled with PCM.

Thermal control of the balloon is not a major concern if a high temperature resistant material is used for the hermetically sealed envelope. Options involve the use of layers of heat resistant silica fabric wrapped around a conventional envelope, but these are less efficient in terms of mass. Elements for the selection of the envelope material are density, mechanical properties at high temperature and the possibility to manufacture thin tight films. The most promising candidate is PBO (Polyphenylene-Benzobis-Oxazole), which gives a safety factor of about ten for temperatures up to 920 K. Although the material is commercially available, technological feasibility has to be verified and a complete characterisation in a simulated Venus environment has to be carried out. Also, the inflation and the characteristics (strength, tightness, etc...) of a densely packed PBO balloon envelope, after the long duration storage in space environment of the several months cruise to Venus has to be demonstrated.

## **2.2 ORBITERS**

In order to place a substantial mass in Venus orbit, without using a very large launch vehicle, two techniques involving interaction with the atmosphere are now available. The first is aerocapture, in which the spacecraft is targeted accurately to the level where atmospheric drag reduced its velocity relative to the planet to the point where it is captured in orbit. Obviously, erroneous targeting, or poor knowledge of the atmospheric conditions, can result in the spacecraft crashing into the planet, or missing it altogether. In either case the margin available is not great. The second, related technique is aerobraking, where the spacecraft achieves initial orbit by conventional means, i.e. using a retro-motor, and the orbit is subsequently adjusted by allowing small amounts of drag to occur near periapsis. By using many orbits over an extended period of time, aerobraking can be gradually increased and the risk is thereby reduced.

### **Orbit Insertion By Aerocapture**

The aerocapture manoeuvre begins with a shallow approach angle to the planet, followed by a descent to relatively dense layers of the atmosphere. Once most of the needed deceleration is reached, the vehicle manoeuvres to exit the atmosphere. To account for the inaccuracies of the atmospheric entering conditions and for the atmospheric uncertainties, the vehicle needs to have guidance and control as well as manoeuvring capabilities. Most of the manoeuvring is done using the lift vector that the vehicle's aerodynamic shape (i.e., lift-to-drag ratio, L/D) provides. Upon exit, the heat shield is jettisoned to minimize heat soak and a short propellant burn is accomplished to raise the orbit periapsis. The entire operation requires the vehicle to operate autonomously while in the planet's atmosphere. Generally, because aerocapture entry velocities are very high, the integrated heat loads are higher than a direct entry and landing. Leading to requirements for new light-weight thermal protection materials for use on the aeroshell.

Aerocapture will be attempted for the first time at Mars, with the Mars Surveyor 2001 orbiter mission. In this case the entry vehicle is a  $70^\circ$  half-angle cone with a 2.4 m diameter base, providing an L/D of 0.18 at an angle of attack of  $11.1^\circ$ . Several aspects make aerocapture at Mars more challenging than at Venus. On Mars, random variations in density can occur rapidly due to the effects of a dust storm, and Mars rotates much faster than Venus and has gravitational "bumps" that create time-varying perturbations.

### Orbit Optimization By Aerobraking

Aerobraking has been successfully implemented on the Magellan mission to Venus. In this case, the aerobraking surfaces were the spacecraft itself and the solar arrays. No special coating or thermal protection systems were added to the spacecraft, although the spacecraft was configured before atmospheric entry to have an aerodynamically stable shape. The orbit of

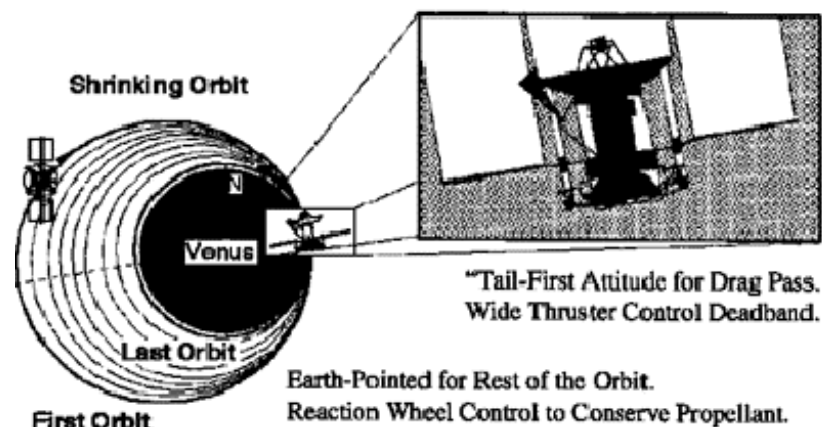


Figure 3. Aerobraking At Venus With Magellan

Magellan was circularized during a 70 day aerobraking phase which ended in August 1993. The orbit apoapsis was lowered from an altitude of 8400 km down to 541 km, providing a  $\Delta V$  of about 1.2 km/s. This manoeuvre would have required an order of magnitude more propellant than Magellan was carrying at the time if done with a chemical burn. The surface temperatures measured on the solar panels did not exceed  $84^\circ\text{C}$ . Several small manoeuvres had to be executed to keep the atmospheric trajectory of the spacecraft within the acceptable corridor.

## 3 GLOBAL CHARACTERISTICS OF VENUS

As a prelude to the more detailed discussion of the Venusian atmospheric and surface environment, we present in this section a brief review of the main planetary parameters and characteristics of Venus and the other terrestrial planets.

### 3.1 INTRODUCTION

Venus is the nearest planet to the Earth, both in terms of distance and in overall size and mass (Table 1). However, Venus has lost most of its atmospheric and surface water, probably as an indirect consequence of its greater proximity to the Sun. A further result is that much of the carbon dioxide which on the Earth has been processed by the oceans to produce carbonates is still free in the atmosphere on Venus. This in turn means that Venus' atmosphere is very massive by terrestrial standards, with a surface pressure of almost 100 bars and a searingly high surface temperature in excess of 730 K. The basic physical characteristics of the planet are summarised in the following table, where the second column shows the ratio with the Earth:

	Venus	$\frac{\text{Venus}}{\text{Earth}}$
Mass	$4.87 \times 10^{27} \text{ g}$	0.815
Radius	6051 km	0.95
Density	$5.2 \text{ g/cm}^3$	0.943
Equatorial gravity	$887 \text{ cm/s}^2$	0.907
Rotation period	243 days (retrograde)	243
Orbital period	224.7 days	0.615
Obliquity	$2.7^\circ$	0.115
Orbital eccentricity	0.007	0.407
Mean distance from Sun	$108.2 \times 10^6 \text{ km}$	0.723

Table 1. Basic data for Venus, and the ratio to the value of the same parameter for the Earth

The rotational axis of Venus is nearly perpendicular to the ecliptic, and the orbit is nearly circular, so seasonal changes in the climate are probably very small. The Venusian year is 224.7 Earth days, longer than the time for Venus to rotate on its axis, which is 243 Earth days. The solar day, defined as the time for the Sun to go from noon to noon as seen from the surface of Venus, is about 117 Earth days. This very slow rotation of the solid body of Venus is retrograde, i.e. backwards compared to the other planets, a curious state of affairs which is difficult to explain in terms of presently accepted models of the evolution of the Solar System (q.v.). The resulting absence of a strong Coriolis force, near the surface at least, is a major factor in determining the structure and dynamics of the atmosphere. Persistent high winds of the order of 100 metres per second are observed near the cloud tops, 50 or 60 km above the surface, where the density of the air is similar to that near the ground on the Earth. It is not known reliably how these are produced.

The surface of Venus is obscured at visible wavelengths by planet-wide cloud cover, the upper layers at least consisting of sulphate aerosol similar to the much thinner layers of volcanic origin found in the terrestrial stratosphere. On Venus, the clouds have complex layered structure and, although never completely absent, very variable coverage. They play an even larger rôle in the energy balance of the planet than clouds do on Earth, through their contribution to the atmospheric ‘greenhouse effect’.

Localised dynamical or ‘weather’ activity on Venus, so far as is known at present, is dominated by four main phenomena: the cloud-top zonal super-rotation, the ultraviolet markings and associated planetary waves, cumulus dynamics in the deeper layers, and the double vortex structures at the poles. These and other types of activity, such as atmospheric tides, are no doubt linked to each other and to the general circulation, but all remain poorly understood (see §5).

### 3.2 Energy Balance

At a distance of 108.2 million km, Venus is closer to the Sun than the Earth by a factor of about  $\sqrt{2}$ , and so has about twice the incidence of solar energy. It is also much hotter at the surface, nearly  $2\frac{1}{2}$  times more than the terrestrial mean of about 300 K. These facts are not simple to reconcile, however, because the ubiquitous and highly reflective cloud cover on Venus reflects 76% of the incoming solar flux and this results in a *smaller* net solar constant

for Venus than for Earth. The high surface temperature must, therefore, be due to 'greenhouse' warming produced by the thick, cloudy atmosphere, possibly augmented by a contribution from the internal heat of the planet.

It is not simple to prove that the observed atmospheric conditions can in fact generate such a large 'greenhouse' effect. The problem is that the massive amounts of carbon dioxide are very effective at blocking the emission of thermal infrared radiation, but only at those wavelengths where the gas has absorption bands, which are far from covering the entire spectrum. Moderate amounts of water vapour are also required, and even then considerable spectral gaps or 'windows' remain. These could be blocked by the clouds, since liquid or solid absorbers present some opacity at every wavelength, the details depending on composition and particle size. The problem for early theorists was that using clouds to 'close' the greenhouse also tended to block the incoming sunlight, so that the calculated equilibrium temperature of the surface remained well below that observed.

This problem began to be resolved when it was realised that the clouds are made of sulphuric acid droplets, at least in the higher, most easily measured layers. These have the property of being highly absorbing at thermal infrared wavelengths, while being nearly conservative scatterers in the visible and near infrared. Thus, the clouds tend to diffuse downwards those of the incoming solar photon that they do not reflect to space, while blocking thermal emission from the lower atmosphere and surface. This explains the result, surprising at the time, that the Venera landers in the 1970s were able to photograph the surface in natural light. It also means that radiative transfer models, involving weak as well as strong bands of CO<sub>2</sub> and H<sub>2</sub>O, plus those of the minor constituents CO, HCl and SO<sub>2</sub>, can account for the high surface temperatures by careful incorporation of the scattering and absorbing properties of the clouds.

The total solar energy diffusing through the cloud cover on Venus corresponds to about 17 watts per cm<sup>2</sup> of surface insolation on the average, about 12% of the total absorbed by the planet and the atmosphere. The high opacity of the gaseous atmosphere and cloud at longer wavelengths requires the surface to reach temperatures high enough to melt zinc before the upwelling flux is intense enough, and at shorter wavelengths, so that equilibrium is attained. An airless body with the same albedo and at the same distance from the Sun as Venus would reach equilibrium for a mean surface temperature of only about 230 K. This 500K greenhouse enhancement of the surface temperature compares with only about 30K on Earth and 10K on Mars.

### 3.3 Composition.

The primordial atmosphere of Venus which originally formed with the solid body, like those of the other terrestrial planets, was likely to have been lost in the distant past as the young Sun went through phases of high activity. The present atmosphere would have been produced much later by outgassing from the crust, a process which we observe today as active volcanism, and by the influx of cometary and meteoritic material, which is also still going on. The relative contributions of these distinct sources can, to some extent, be deduced from the data which is gradually being accrued on the composition, and in particular the isotopic ratios, in the contemporary terrestrial planet atmospheres, and in comets and meteorites.

	<u>Venus</u>	<u>Earth</u>	<u>Mars</u>
Carbon dioxide	.96	.0003	.95
Nitrogen	.035	.770	.027

Argon	.00007	.0093	.016
Water vapour	~.0001(?)	~.01	~.0003
Oxygen	.0013	.21	~ 0
Sulphur dioxide	.00015	.2 ppb	
Carbon monoxide	.00004	.12 ppm	.0007
Neon	5 ppm	18 ppm	2.5 ppm

Table 2. Composition of the Terrestrial Planet Atmospheres, as fractional abundances except where ppm (parts per million) is stated.

Comets are a rich source of volatile compounds such as carbon dioxide, water vapour, methane, and ammonia. If the last of these was the source of the nitrogen now present, and we allow for processes such as the production of argon by the decay of radioactive potassium in the crust, the contemporary atmosphere could all be of external origin. On the other hand, the high abundance of sulphur in Venus' clouds is strongly suggestive of extensive volcanic activity and volcanoes are also prolific sources of carbon dioxide, nitrogen and the other gases required to explain present-day Venus.

Apart from carbon dioxide and water vapour, Venus' atmosphere consists primarily of inert gases, particularly nitrogen and argon (Table 2). The amount of water present as gas and bound up with sulphuric acid and other compounds in the clouds is between ten and one hundred thousand times less than exists in the oceans and atmosphere of the Earth. Thus Venus is overall very dry compared to the Earth while, at the same time, deuterium is about one hundred times more abundant on Venus than Earth. This suggests that Venus had much more water initially, but that most of it has been lost. Loss takes place by dissociation of the water in the upper atmosphere by solar ultraviolet radiation, and the subsequent escape of the hydrogen. Both deuterium and normal hydrogen escape from the atmosphere while there is free water on the surface, but the heavier isotope escapes less efficiently, leading to the observed fractionation.

The loss rate of the water depends strongly on its abundance in the relatively cool middle atmosphere, as well as the intensity of the solar ultraviolet flux. Models of the process suggest that Venus could have lost an ocean of present-day terrestrial proportions in only a few hundred million years. The oxygen produced at the same time is too massive to escape at any significant rate, according to Jeans' formula, and must remain on the planet, presumably most of it bound chemically within the crust, mainly as carbonate rocks. As on the Earth, this process would remove atmospheric carbon dioxide efficiently so long as liquid water was available. Once the free water was all used up, the mixing ratio of water vapour in the upper atmosphere would fall sharply and the loss rates of both forms of hydrogen, and the take up of oxygen into minerals, have also declined to the present relatively low levels.

In the modern atmosphere of Venus, chemical reactions coupled with transport and radiative processes regulate the abundances of the most important minor constituents. The most important are the cycles involving water vapour, sulphuric acid and their products, which maintain the cloud layers, and which probably also involve reactions between the atmosphere and the surface. Another is that which gives rise to the observed distribution of carbon monoxide. CO is very abundant (mixing ratios of the order of a few parts per thousand by volume) in the upper atmosphere of Venus, as would be expected from the action of solar ultraviolet radiation on carbon dioxide. It is strongly depleted in the cloud layers (<1 ppmv), again not too surprisingly, since it is involved in reactions with SO<sub>2</sub> and the other species which make up the sulphur cycle. Below the clouds, and near the surface, however, the

carbon monoxide value recovers to around 30 ppmv, and shows a marked equator-to-pole gradient. It seems likely that CO is transported rapidly down from the thermosphere in the polar vortices (see below), to the troposphere where it is gradually removed by reactions in the hot lower atmosphere and at the surface.

## 4 SURFACE PROPERTIES

In this section the surface of Venus is reviewed, based mainly on the radar data from Magellan and Pioneer. The influence of the surface on the atmosphere is essentially unknown; the topography is less extreme overall than on Earth or Mars, but still substantial and likely to produce waves in the dense, slowly-moving lower atmosphere which propagate and grow with height. There is almost certainly a large chemical coupling between surface and atmosphere, as well. In addition to the strong possibility of active volcanism, which probably plays a role in maintaining the cloud layers, the high temperatures and corrosive minor constituents in the atmosphere suggest that chemistry must be a major factor in controlling the surface environment.

### 4.1 TOPOGRAPHY

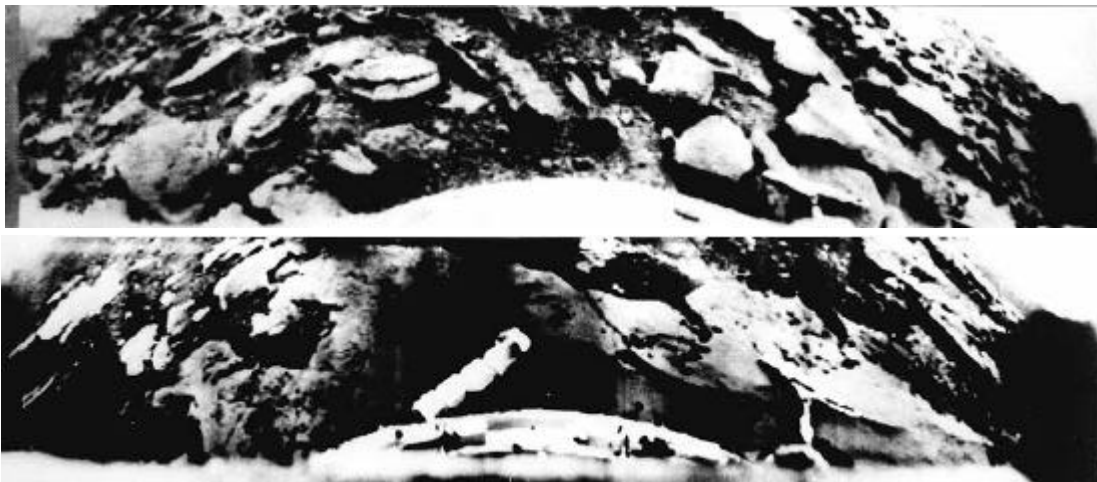


Figure 4. Venera 9 and 10 views of the surface of Venus.

The only visible images of the surface of Venus came from the Venera and Vega landers, and were obtained in natural light, showing the ability of a few percent of the sunlight incident upon Venus to diffuse through the clouds. The surface of Venus is revealed as a sterile, scorched desert, dominated by the boulders which appear strewn about the landscape. However, there are significant differences between, for example, the Venera 9 and 10 landscapes, which are located about 2500 kilometres apart in the equatorial regions of Venus. The Venera 10 panorama shows a flat, rocky plain with a low outcropping of rock. It seems likely that Venera 9 landed on a slope (sensors on the spacecraft showed that it was tilted at about  $30^\circ$  to the horizontal after settling on the surface) and that the boulders are 'scree' or rubble from the break-up of the faces of the hill upon which the spacecraft sits. Venera 10, on the other hand, sits on a fairly flat, rocky plain.

Both areas contain features which are consistent with the ejecta and lava flows associated with volcanic activity. The stones lying on the ground look much like an alluvial deposit, some of them having dark bands and others a patchy appearance. They rest on a mottled deposit which looks like crushed basaltic rock. The black regions in this background are too dark to be shadows or even dark minerals; they must be depressions or fissures between the rocks. Such an interpretation is re-enforced by the very obvious sharp edges on some of the boulders, clearly the result of fracturing by some geological process. The fracturing could

have been very ancient, since running water, large daily or seasonal temperature changes and wind erosion are not available to weather rocks on Venus as they do on Earth. On the other hand, some of the rocks do show evidence of erosion. Understanding the nature of this process is one of the major unsolved problems raised by the Venera pictures. Chemical erosion by acidic vapours in the atmosphere, or melting of volatile components of the rocks, are candidates under consideration. Such processes certainly must occur, but they probably lead to much slower weathering than takes place on Earth.

The level of illumination in the Venera pictures was higher than had been expected. Even with the Sun  $60^\circ$  above the horizon, it was thought that the thick clouds would prevent more than a trace of sunlight from reaching the ground; instead a light level which has been compared to that on Earth during a thunderstorm was discovered and the searchlights which the spacecraft carried were unnecessary. In 1978, radiometers on the Pioneer Venus probes measured the solar flux and found that 2% of the total falling on the planet actually reaches the surface without being absorbed. This is a large fraction considering the thickness of the cloud layers and shows, as do other measurements, how reflective the cloud droplets are on the whole. They diffuse the radiation thoroughly by scattering each photon dozens of times during its passage through the atmosphere, but do not absorb as strongly as terrestrial clouds would.

Other instruments on the Veneras, principally gamma-ray spectrometers, revealed the chemical composition of the surface rocks. At both sites, the abundances of the naturally radioactive elements uranium, thorium and potassium were consistent with a composition like terrestrial basalt. The density measurements, by gamma-ray backscattering, of  $2.7\text{--}2.9\text{ g/cm}^3$ , are consistent with this conclusion. Thus it seems that Venus formed in a manner similar to the Moon, Earth and Mars; by condensing from a molten protoplanet into shells, the outermost of which (the crust) is primarily composed of the most fusible material, the basalts.

Radar mapping by the Pioneer Venus Orbiter in 1978 slowly built up maps strip by strip as the orbit precessed, taking a Venus year of 243 days to complete. The radar map revealed a planet that is characterised primarily by smoothly rolling plains (about 70% of the surface area), with clearly distinguished highland (about 10%) and lowland (about 20%) regions. There are four principal highland areas of Venus, thought to correspond to areas of mantle upwelling, and which are roughly comparable to the continents on Earth. They are (1) Ishtar Terra, (2) Lada Terra, (3) Aphrodite Terra and (4) Beta, Phoebe and Themis regiones. In addition, there is Lakshmi Planum, a 3-4 km high plateau, bordered by mountainous ridges.



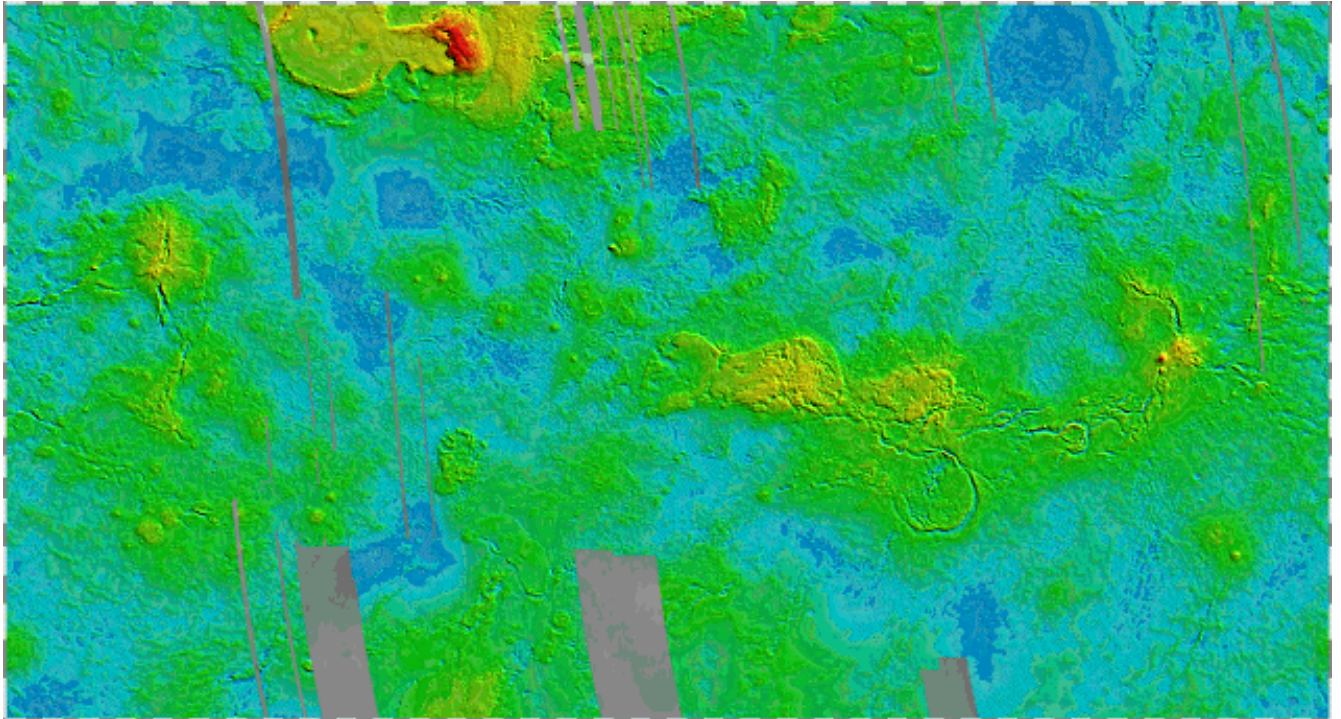


Figure 5. Magellan radar map of the surface of Venus.

Ishtar covers an area comparable to Australia and rises steeply from the surrounding plains at about  $70^{\circ}\text{N}$ . The western part is a high plateau (3 km above the mean radius of Venus) bordered by tall mountains that reach a further 3 km in altitude. In the middle of Ishtar stand the *Maxwell Montes*, mountains that, in rising to 11 km above the mean, would tower above Everest, and are much more steeply-sided. The existence of such a steep and massive mountainous feature on Venus speaks of some unique feature deep within the crust of the planet, which created the mountain, and continues to support its massive bulk. The prime candidates are a particularly energetic collision between surface plates, trying to move sideways into each other (in which case Maxwell would be analogous to the Himalayas on Earth) or very vigorous (by the standards of molten or plastic rocky material) upwelling in a large 'hot spot' in the crust.

Stretching for about 10000 km along and south of the equator, *Aphrodite Terra* is the other very prominent highland region, in this case one that covers an area about equal to that of Africa, but with a more elongated shape which has been likened to a scorpion. The western end of Aphrodite is made up of two elevated, ancient, fractured plateaux. The highlands to the east of these extend for 5,000 km and contain steep valleys, some of them on an enormous scale not found on Earth. The largest, Diana Chasma, is on a similar scale to the vast Valles Marineris on Mars.

The lowlands on Venus, generally designated *Planitiae*, are generally the most featureless regions in terms of tectonic and volcanic structures, probably a result of having been flooded by lava relatively recently.

The most recent mission to Venus, Magellan, was dedicated to radar imaging of the surface and the study of these features in more detail, with a resolution of 75 to 120 m. Magellan resulted in an explosion of knowledge about the surface of Venus, taking it from the least to one of the best explored terrains in the Solar System. The individual features seen in the Magellan radar pictures are generally of three types: volcanic, tectonic or impact features. The first of these refers not only to the volcanoes themselves, but also to vast lava flows which are seen filling the plains which cover huge areas of the planet. Some areas appear to

have been flooded several times, with fresh flows partially covering the earlier ones. Because of this resurfacing, much of the surface of Venus is geologically very young. There are relatively few impact craters, compared to the Moon and Mercury, for example, and those that are seen are uniformly distributed over Venus, suggesting that the renewal of the surface has been planet-wide.

### **Volcanic Features**

About 90% of the surface of Venus is made up of features attributable to volcanic activity. The volcanoes themselves can be divided into three groups according to their size. The largest are the shield volcanoes, of which 156 have been identified in the Magellan maps, mostly in high regions lying 3-5 km above the surrounding area. They have lava flows which often extend for hundreds of kilometres, radiating away from a central crater or *caldera*.

The intermediate-sized volcanoes are further sub-divided into *anemones*, *ticks*, and *arachnids*, depending on their appearance. Anemones are relatively rare, with only 25 so far identified. They are characterised by 'hairy' flow patterns typically 50 km across, radiating out from a central source of magma. About 50 *ticks* have been found on Venus. These are flat, circular volcanic domes about 25 km in diameter flanked by strongly defined radial ridges and troughs, giving them an appearance reminiscent of the eponymous insect. The arachnoids may be originally volcanic mounds, which have collapsed, cracking the crust and producing a distinctive insect-like shape.

Similar to ticks but without the 'legs' are the steep-sided, flat-topped volcanoes known as *pancake* domes. These have well-defined circular outlines and patterns of radial fractures, especially at the edge, and one or more small calderas near the centre. Over 150 of this type have been recognised, and they are often found in groups, sometimes overlapping one another. The pancake shape probably means that the lava which formed them was of higher viscosity than that emanating from the more earth-like large volcanoes, and so did not flow as freely or as far. Such evidence as exists, from the Venera landers in particular, for the composition of the plains on Venus suggests that they are made of basaltic material. Were the pancake domes and related features made of something more akin to granite, this might have the properties required to explain the observed formations. The nearest analogues to pancake domes found on the Earth are on the sea bed, perhaps suggesting that the surrounding high-density fluid also plays a key role in the cooling and solidification of the dome.

At the other end of the viscosity range are the lavas which form remarkable rivers on Venus, some of them extending thousands of kilometres from their volcanic sources to the lava-filled flood plains. Many of these appear to have eroded deep valleys, like rivers on Earth, which suggests that they have flowed consistently for long periods of time. Obviously, water is not involved in the case of Venus, but rather something which has a melting point which is not too different from the mean surface temperature on Venus. Assuming this temperature has not changed too much since the rivers flowed, the most likely candidate would seem to be a carbon-rich mineral such as carbonatite, although there are others that cannot be ruled out, even including the low melting-point metals like lead or tin. Although the plains are clearly solidified now, we cannot know for certain that these rivers, if they are different in composition from the large-scale flows, are not still running in places.

In many places on Venus, volcanoes are clustered together to form *shield fields*, which contain large numbers of small domes or vents, each typically a few km across but collectively covering an area of more than 10,000 square kilometres. There are hundreds of these on Venus, some with extensive lava flows surrounding them, while others are located within tectonic structures, as described below. Finally, around a hundred volcanic calderas have been identified on Venus which are apparently sources of lava flows but not associated with cones or domes.

## Tectonic Features

Tectonic features are those produced by movements of the crust. Unlike its near-twin, the Earth, Venus does not seem to have continent-sized plates which move relative to each other, causing continental drift. The reason is probably that some combination of the composition of the crust (such as its likely low water content), its thickness, and the high surface temperature, weakens the rock on Venus to the point that it cracks more easily than on Earth, preventing the formation of major plates. The extensive network of narrow linear structures which is seen covering the surface of Venus could be evidence for this, if they show where the crust has been cracked by the same kind of crustal flexing that produces plate tectonics on Earth.

The tectonic features are given Latin names and most fall into the categories of *tesserae*, *chasmas*, or *montes*. The *tesserae* (the word means 'tiled') are networks of intersecting ridges and troughs which appear to have formed when plates were forced together by movement of the crust. The fact that they are generally overlaid at their flanks by other types of terrain suggest that *tesserae* represent the oldest regions still apparent on the Venusian surface. A *chasma*, as the name suggests, is a chasm or broad trench, or sometimes several of these running in parallel, and bounded by cliffs. They are produced when sections of the crust pull apart. The *montes* of course are mountains or mountainous provinces, the outstanding example of which, Maxwell Montes, is remarkably high, massive, and steep, as already noted.

Another category of features common on Venus are the *coronae*. These are large, circular patterns of ridges and troughs ranging in diameter from 75 to over 2000 kilometres. Although basically tectonic features, the *coronae* have their origins in volcanic activity below the surface, the crust being bulged and cracked by convection in the sub-surface lava field, which may or may not break through the surface. There are various sub-species of *coronae*, including some which may have been produced as a result of lava escaping and forming a mound, which then collapses, resulting in further cracking and movement. The local fracturing and movement is greatest in regions where the crust is particularly thin or weak, as evidenced by the cracks or *graben* which radiate out of some *coronae*, often extending well outside the region of lava flow. These well-developed examples are the *arachnoids* mentioned earlier. If the 'legs' are conspicuous but the 'body' is underdeveloped then they are termed *novae*, the assumption being that they represent the early stages of corona formation.

## Impact Craters

It has already been noted that Venus has relatively few impact craters compared with Mercury, Mars or the Moon, and that their spatial distribution is uniform and random. A total of 935 have been found in the Magellan maps, and many of the bigger ones have been named after famous women. The largest, Mead, is 280 km in diameter. The marked absence of craters smaller than a kilometre or so in diameter is consistent with the considerable thickness of the atmosphere, which prevents smaller meteors from reaching the surface.

It seems certain that there have been many more crater-forming impacts on Venus than the current record suggests, but that old craters have been obliterated by atmospheric erosion, tectonics and lava flows. It is quite difficult to assess the relative importance of these. Lava flows are the fastest way to resurface the planet, and as already noted there is plenty of evidence for extensive flooding of the plains on Venus. Indeed, this may still be going on to some extent, either continuously or episodically.

Crumbling of the crust, which, as we have seen, occurs on a much smaller scale on Venus than on the Earth, will also remould the surface, especially when it is accompanied by local seepage of lava.

Atmospheric erosion has a major role in obliterating craters on the Earth, but does not seem to be so important on Venus, because those features which are present in the Magellan images

appear to be remarkably pristine. For one thing, the surface winds on Venus are much less than on Earth, and this is probably a more important factor than the high density. Perhaps surprisingly, the action of water, and vegetation, on the Earth seems to be a more effective force for corrosion of the topography than the sulphurous and other compounds which abound in the hot air near the surface on Venus. The higher features on Venus seem to be not only relatively undefiled, but actually plated with something, which gives everything more than about 2 kilometres above the mean surface height an unnaturally high radar brightness. An intelligent guess as to what this is can be made by looking for something which condenses at the temperature and pressure of the observed lower altitude boundary. Results include the element tellurium or, even more intriguingly, iron pyrites, also known as fool's gold from its resemblance to the precious metal.

Impact craters are classified by their appearance into one of six categories. *Structureless* craters are the simplest, and most often the smallest, with a flat and featureless floor. *Central peak* craters have a central uplift that rises above the crater floor and terraced walls. They are particularly circular in outline and range in size from 8 to 79 km, with most between 20 and 30 km. *Double-ring* craters are typically greater than 40 km, with an outer rim and a circular arrangement of inner peaks and ridges. The largest craters on Venus, ranging from 86 to 280 km in diameter, are of the *Multiple-ring* type, with two or more sets of concentric outer walls. *Irregular* craters have non-circular rim outlines and structural disruptions to otherwise flat crater floors, are very common (almost 1/3 of the craters on Venus are of this type) and quite small (less than 16 km across). *Multiple* craters form when a falling body fragments into pieces, each creating a separate impact crater up to 44 km in diameter and sometimes forming a neat chain.

#### 4.2 Magellan Datasets

The Magellan data set on the surface of Venus is available from the National Space Science Data Center (NSSDC) at NASA's Goddard Space Flight Center, the principal archive and distribution centre for all NASA missions. NSSDC has the altimeter/radiometer data collected by Magellan in digital form on CD-ROM and in photographic form. It will have new maps of Venus as they are produced by USGS based on Magellan's data. It has all of Magellan's standard mosaic image products which have been released by the Project in both photographic and digital form on CD-ROMs. It also has press-released images, videotapes, software with which to display the CD-ROM digital images, planning maps of Venus, a fact sheet, and documentation.

The Magellan dataset consists of over 130 CD-ROM disks of digital radar image mosaics and 15 disks of altimetry and radiometry data. These cover most of the first two 243-day mapping cycles around Venus. Magellan has also collected data not involving the radar instrument.

These include:

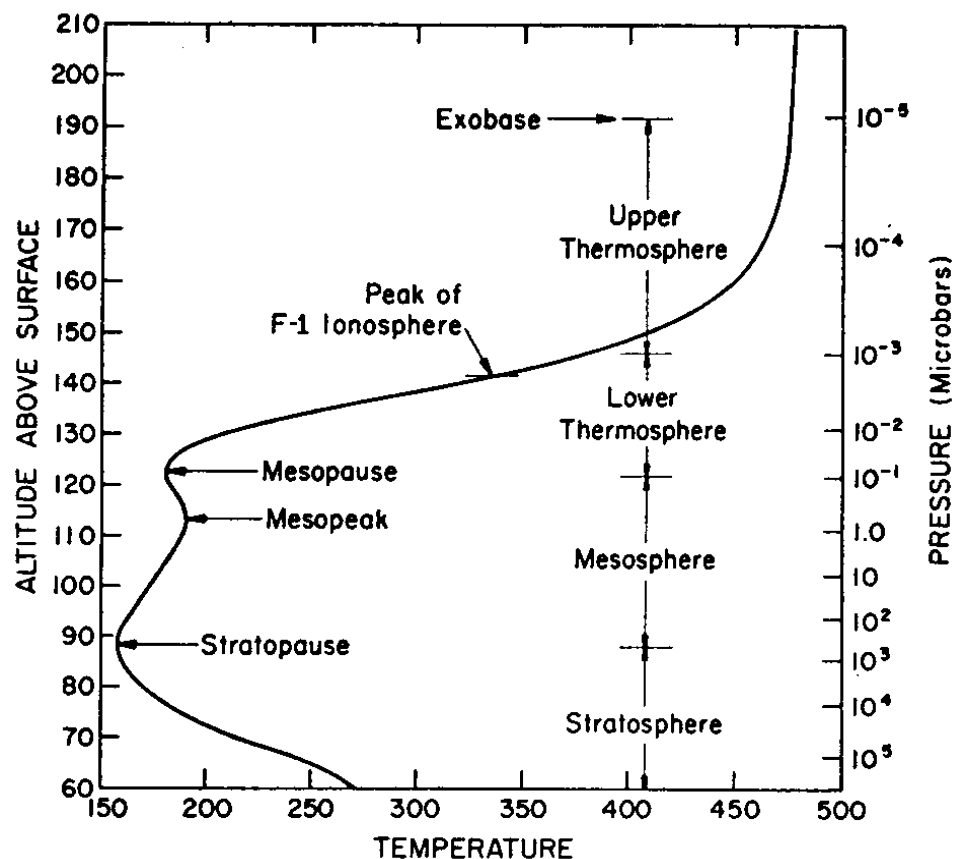
- \* Radio occultation data measuring the effect of Venus's atmosphere on the spacecraft's downlink signal;
- \* Phase, amplitude, and polarization data measuring the effect of the Sun's outer plasma on the downlink signal; and
- \* Drag and torque data during periapsis passages measuring the density of Venus's very high atmosphere.

## 5 ATMOSPHERIC PROPERTIES

In this section we review current knowledge of various aspects of the Venusian atmosphere, including its vertical temperature structure, circulation, and various physical processes such as radiative transfer, boundary layer effects and clouds.

### 5.1 Thermal structure

The solar radiation which penetrates the clouds warms the lower atmosphere, which is prevented by the opacity of the overlying layers from cooling by radiation to space. It therefore forms a deep convective region, the troposphere (Fig. 6). This links the high surface temperature of around 730K, produced by the greenhouse effect as described above, with the level at which the temperature is close to the effective bolometric temperature of Venus (about 230K), where strong radiative cooling to space can occur. The adiabatic lapse rate, which applies when the vertical gradient is controlled by convection, is  $-g/c_p$  or about 10 K km<sup>-1</sup> for Venus. The troposphere, up to about 60 km, has a larger vertical extent than on Earth or Mars, since the larger density and optical thickness produce a much deeper convection zone.

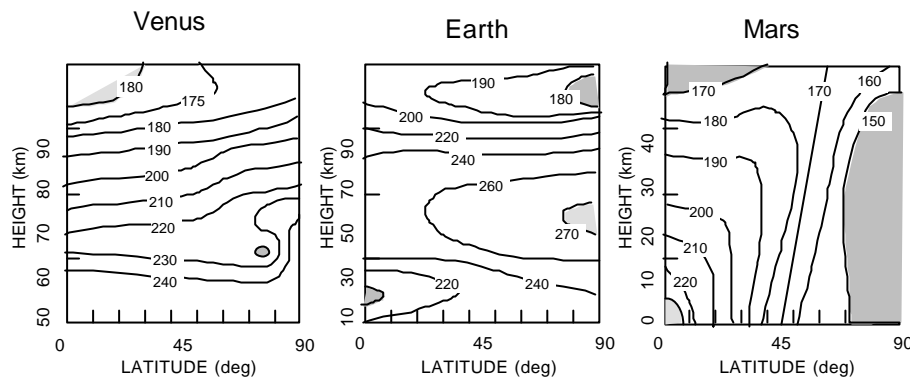


**Figure 6.** Model temperature profiles for the atmosphere of Venus.

Above the troposphere and the clouds top, from 60 to about 100 km, the temperature still maintains a constant adiabatic lapse rate, i.e., it is largely convectively driven, but in the upper part of this lower mesosphere region it tends to become constant with height as the atmosphere becomes optically thinner and, to a first approximation, each layer tends to find a radiative equilibrium temperature. This temperature is determined by the balance between the absorption of upwelling infrared radiation from lower atmospheric layers and from the clouds top and emission to upper layers and cooling to space, with no significant absorption of direct solar energy taking place until higher levels.

On the Earth, this situation is altered by the ozone layer, which is responsible for substantial heating which divides the middle atmosphere into the stratosphere, where temperature rises with height, and the mesosphere, where it falls. There is no corresponding effect on Venus (Fig.7), except for small amounts of absorption of solar and thermal energy in the near infrared bands of water vapour and carbon dioxide. This absorption produces a mesopeak (or local temperature maximum) in global radiative equilibrium models around 110 km. The

thermosphere is reached at heights of more than 120 km above the surface. Here the atmosphere is very rarified, and short-wavelength solar radiation and energetic particles drive rapid photochemical reactions and produce high temperatures by day. At night, when the heating is absent, efficient radiative cooling by CO<sub>2</sub> results in a rapid decline of temperature, so there is a sharp gradient across the terminator, from over 300 K on the illuminated hemisphere down to 100 K or less on the night side.



**Figure 7.** Meridional temperature cross-sections for the atmospheres of the terrestrial planets. Schematic only - for original data, see (a) Pioneer Venus Orbiter Infrared Radiometer measurements (Taylor et al., 1981), (b) Nimbus 7 Stratospheric and Mesospheric Sounder data (Taylor, 1987), and (c) Mariner 9 Infrared Interferometer Spectrometer data (Conrath, 1981).

Data on the planet-wide variability of temperature below the clouds is sparse, and is still mainly limited to the small number of direct measurements made by the various entry probes. Remote sensing of the lower atmosphere has recently been shown to be possible, in particular at near infrared wavelengths, as well as in the microwave, but uncertainties in the distribution of opacity sources in both spectral regions make these difficult to interpret in terms of temperature variations. Such data as exist show that the temperature gradients in the lower atmosphere are close to adiabatic in the vertical, and close to zero in the horizontal, as would be expected theoretically from the high opacity and high density.

Above the clouds, where the density is lower, more variability is expected and has been observed. Fig. 8 shows the time-averaged (72-day mean) global temperature field from the Pioneer Orbiter in which several features clearly related to the general circulation stand out. Note the 'polar warming', in which the equator-to-pole temperature gradient leads to higher temperatures over pole than equator, in spite of the fact that the trend in radiative heating is in the opposite direction; and the 'polar collar' feature in the meridional cross-section, an intense ribbon of cold air surrounding the pole at about 65 degrees latitude. These are both features of dynamical origin, related to the zonal super-rotation and the polar vortex respectively. The meridional average temperature field is dominated by the diurnal variation or solar tide, the temperature cycle which is induced by the apparent motion of the Sun overhead. This contains a whole spectrum of Fourier components, because the forcing is non-sinusoidal; the actual atmospheric response depends on the mean wind and the interference between the various components. The solar tide on Venus and classical tidal theory as originally developed for the Earth can be reconciled, provided that a realistic representation of the zonal wind is incorporated.

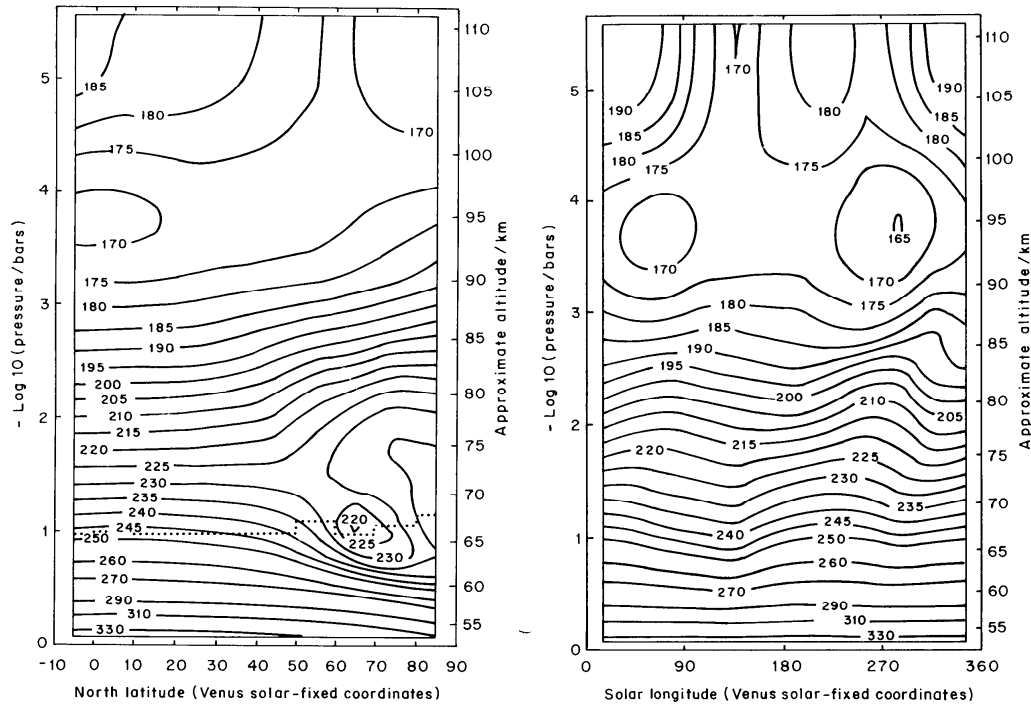


Figure 8. Time averaged temperature fields in the middle atmosphere of Venus (Schofield and Taylor, 1983). (a) The zonal mean field, and (b) the variations around a latitude belt from 0 to 300 N, both plotted against pressure and approximate height. The horizontal stepped line represents the retrieved mean cloud top height.

### 5.2 Clouds.

Venus is completely enshrouded by clouds in a complex layered structure over 30 km deep; their properties are summarised in Table 3. As already noted, the clouds are a key part of a highly interactive climate system. Concentrated  $\text{H}_2\text{SO}_4$  droplets have properties which contribute very significantly to the atmospheric “greenhouse”, by scattering conservatively at short (solar) wavelength radiation while strongly absorbing at long (planetary) wavelengths. Changes in the optical properties or depth of the cloud layers, for example if the clouds dissipated or changed their composition, due to a reduction in the supply of  $\text{SO}_2$  and other source gases from volcanoes, or to an instability in the dynamical regime, would cause the lower atmosphere and surface gradually to cool down (or, conceivably, to become even hotter). Changes in temperature would be likely to further modify the atmospheric circulation and the formation of clouds, producing feedback which seems more likely to be positive (accelerating the change) than negative. At present the radiative, dynamical and chemical processes appear to be in balance, but the stability of the currently observed state may be precarious.

	<u>Venus</u>	<u>Earth</u>	<u>Mars</u>
Fractional coverage	1.00	0.40	0.05 (cloud); 0 - 1.0 (dust)
Typical optical depth	25-40	5-7	0.01-1.0; 0.2-6 (dust)
Composition	$\text{H}_2\text{SO}_4.\text{H}_2\text{O}$	$\text{H}_2\text{O}$	$\text{H}_2\text{O}, \text{CO}_2$ Magnetite etc.(dust)
Number density, liquid	50-300	100 -1000	0

(cm <sup>-3</sup> )			
Number density, solid	10-50	0.1 -50	30 -1000 (near surface)
(cm <sup>-3</sup> )			
Typical mass loading	0.01 -0.1	0.1 -10	.0002 -0.1
(g.m <sup>-3</sup> )			
Main production process	chemistry	condensation	condensation windblown (dust)
Equivalent depth (mm)	0.1 -0.2	0.03 -0.05	1 -100
Effective radius (μm)	2 -4	10	0.4 -2.5
			(dust)
Main forms	stratiform	stratiform, cumulus	stratiform, mixed (dust)
Temporal variability	slight	high	high
Dominant heat exchange process	radiation	latent heat	radiation

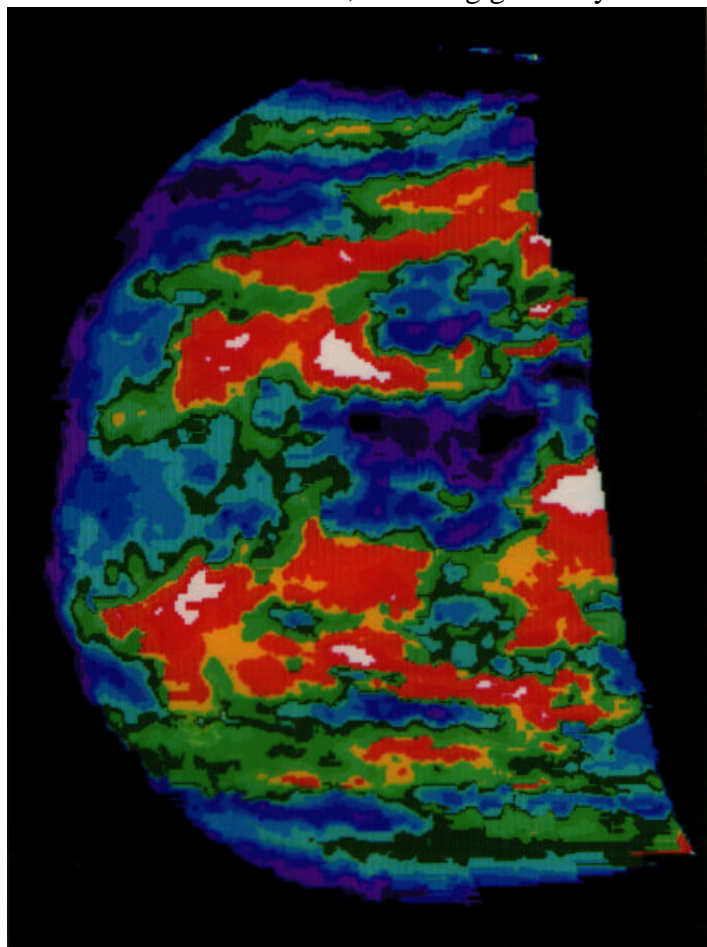
Table 3. Properties of clouds and dust in the terrestrial planet atmospheres.

Most of our detailed knowledge of the cloud properties comes from optical measurements: polarimetry as a function of phase angle from the Earth, nephelometry and particle size measurements from entry probes, and visible, ultraviolet and infrared remote sensing from orbiter and fly-by spacecraft. It has recently been discovered that near-infrared spectroscopic measurements in atmospheric "windows", that is, wavelength regions where the main atmospheric gases are weakly absorbing, penetrate the clouds, in some windows all the way to the surface. This type of observation reveals the cloud morphology all over the planet, and shows that the integrated vertical opacity is very variable, by factors of 20 or more. Thick and relatively thin clouds form patterns suggestive of large-scale cumulus dynamics, presumably with the cloud material actively condensing and dissipating in rising and falling air associated with weather systems (fig.9), although the details are lacking because of a shortage of high resolution data in space and time.



Figure 9. 2.7  $\mu\text{m}$  image of Venus from Galileo-NIMS

At the Pioneer Venus Large Probe entry site (4.4N, 304W), the main cloud deck extends from about 47 to about 67 km above the surface, declining gradually at the upper boundary with a



scale height of a few km, and with a thin haze layer some 13 km deep below. Within this vertical structure, detailed, and presumably variable, layering occurs and particles of different sizes congregate at different height levels. The particles range in diameter from less than 1 to over 30 microns and tend to a trimodal size distribution, with the commonest diameters falling towards the ends of the overall range, and in the 2-3 micron region. It is these intermediate size or 'mode 2' droplets which are visible from outside Venus and for which spectroscopic, polarimetric and other evidence yields a composition of 75%  $\text{H}_2\text{SO}_4$  and 25%  $\text{H}_2\text{O}$ . The composition of the smaller drops, which form an aerosol haze extending throughout the cloud layer, is unknown. Most of the mass of the clouds is in the biggest drops, for which there is some inconclusive evidence of a non-spherical shape, implying a solid composition, perhaps crystalline sulphur. The formation of the cloud droplets can be explained by a model in which  $\text{H}_2\text{O}$  and  $\text{SO}_2$  (the latter possibly of volcanic origin) combine photochemically near the cloud top level. It is more difficult to explain the size distribution, particularly the existence of more than one mode. Compositional contrasts and dynamical effects may be at work but once again the observations which would elucidate these are lacking.

There is some evidence in infrared maps from the Pioneer Venus orbiter of cold, high cloud, presumably condensed  $\text{CO}_2$ , near the tropopause above the dawn terminator, where the atmosphere is coldest as it comes to the end of the long (>50 hours at this altitude) Venus night.

### *5.3 Radiative Transfer In The Upper Atmosphere*

Radiative transfer models for Mars can also be applied to the upper atmosphere of Venus, with only a few modifications, because both have a nearly pure CO<sub>2</sub> atmosphere. In the lower atmosphere of Venus, however, the prevailing temperatures and pressures are much higher and the codes have to be modified and extended to take this extended regime into account.

Less work has been done on radiative transfer models for Venus in recent years, partly because the studies are generally stimulated by plans for space experiments intended to make relevant measurements, and Mars is now the centre of attention in this respect. A decade or two ago, the situation was reversed, and a great deal of what is now known about radiative conditions on Venus still depends on pioneering work done then, especially by Dickinson. His approach was first to calculate a global mean radiative equilibrium temperature profile, including all processes thought to be of importance for the global heat budget between 65 and 135 km, including non-LTE, with a simple treatment of other aspects like overlapping between lines. (LTE stands for local thermodynamic equilibrium, the condition which applies in general when the populations of molecular energy levels are determined by collisions rather than radiative exchange). He pointed out the particular importance of including the combination and, especially, the isotopic bands of CO<sub>2</sub>. As for the case of Mars, these weaker bands have a disproportionate effect because radiative transfer of energy can take place over much longer distances than in the strong fundamentals of the commonest isotope. Only bands with a strength at least five orders of magnitude less than the strongest bands can be neglected. This effect is illustrated in Fig.10, which shows the contribution of a single spectral line to the cooling rate from the 10 $\mu$ m level. Even at this low pressure, the core of a strong line is completely absorbing and cooling to space is negligible, except in the line wings. A weak line, on the other hand, emits strongly to space from its Doppler core.

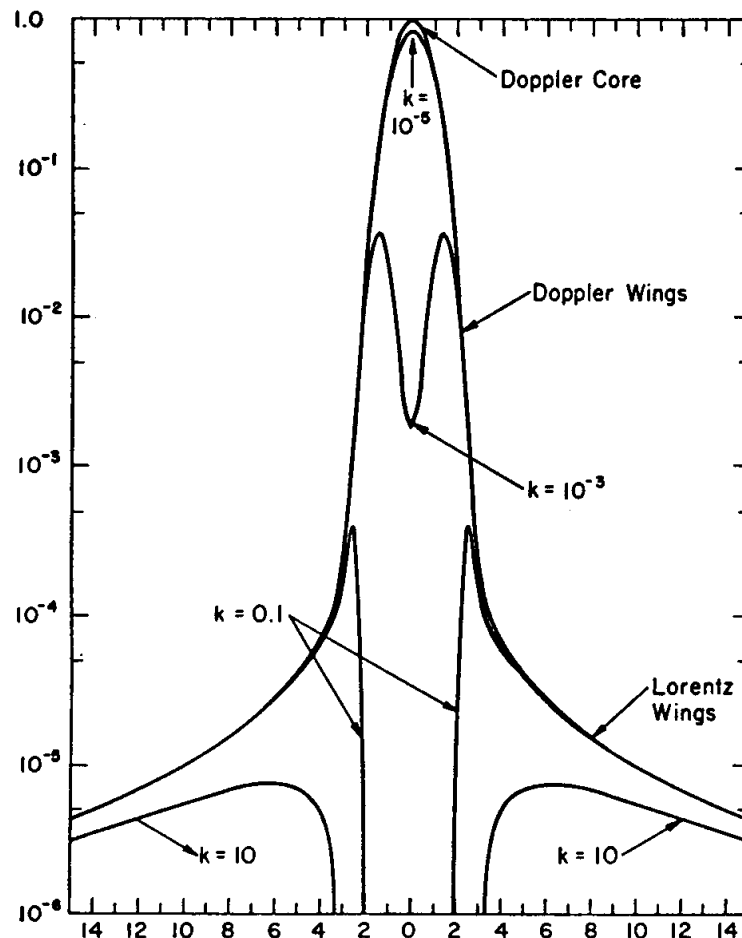


Figure 10. The outer line is the emission profile of a  $15\mu\text{m}$  band line at the  $10\mu\text{bar}$  level on Venus. The other lines are the angle--integrated contributions to the net cooling from that level of a line with the strength  $k$  indicated. The abscissa is in units of the Doppler width (after Dickinson, (1976).

Dickinson refined some aspects of his original work in later revisions, like the collisional coupling in his non-LTE model. For more than a decade his radiative studies have remained unmodified, and guided the work of the group at NCAR and Arizona University that developed the most complete thermospheric global circulation model of Venus to date. In particular he studied the susceptibility to non-LTE of temperature sounding experiments which were at that time (1975) in preparation for space missions to Venus. He showed that the  $15\mu\text{m}$  fundamental band could be considered to be in LTE up to pressures as low as  $1\mu\text{bar}$ , but that at lower altitudes most of the photons escaping to space originated in the hot and isotopic bands. This meant basically that 'wide-band' measurements centred in the fundamental could sound the Venusian mesosphere without taking account of non-LTE, but the use of pressure modulators to sound the lower thermosphere would require a more complex interpretation. Only recently new and more complete non-LTE models developed at the Instituto de Astrofísica de Andalucía in Granada have been applied to the Mars and Venus atmospheres. They improved upon previous codes like Dickinson's in the incorporation of the most recent rate coefficients and spectroscopic data, and by extending the number of  $\text{CO}_2$  levels and transitions studied simultaneously. These models represent useful tools to help in the interpretation of satellite data and the design of radiative parameterizations for global circulation models.

## Cooling rates

All previous studies which computed the cooling and heating rates for the Venusian and Martian atmospheres stress the important role of atomic oxygen: the thermal collision between CO<sub>2</sub> and O is a key process in the cooling of the upper layers of a CO<sub>2</sub> atmosphere.

A typical cooling rate profile produced by the CO<sub>2</sub> 15 $\mu$ m emission on Venus is shown in Fig.11. The maximum cooling of about 130K/day occurs between 120 and 130km and shows a double peak, caused by the fact that the contributions of the main and minor isotopes peak at different altitudes. This large radiative cooling keeps the thermospheric temperature well below Earth's values (300 K versus 1000-2000 K). Again, the emission from the minor isotopes would be much less were it not for the large energy transfer by V-V coupling between their (0,11,0) levels and that of 626.

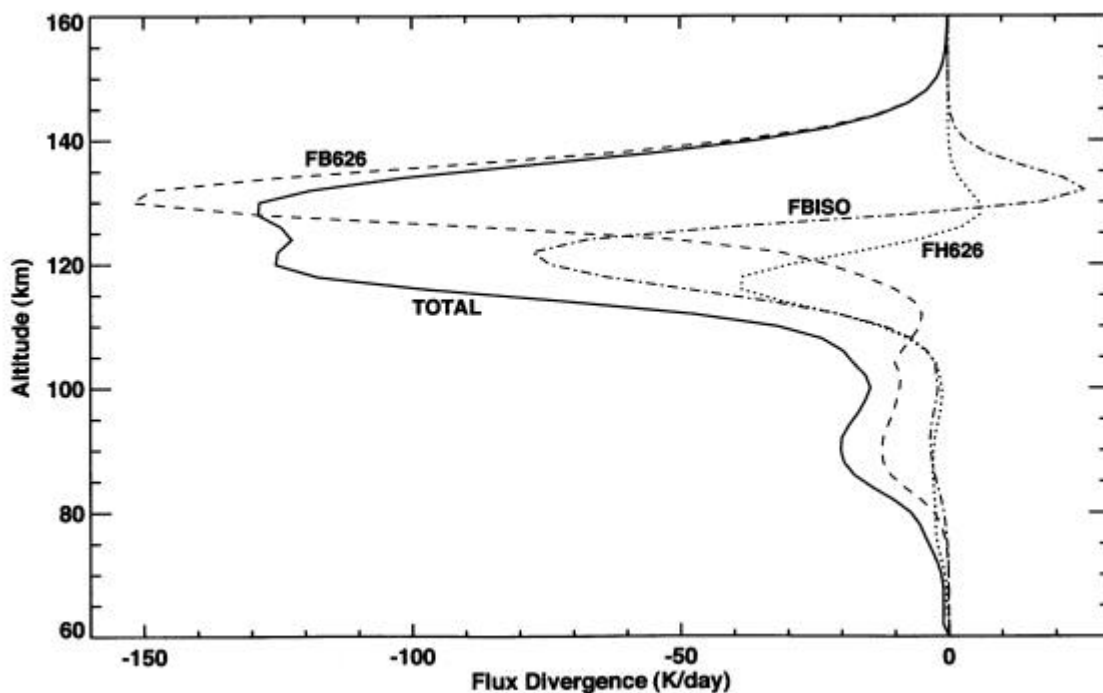


Figure 11. Contributions and total cooling rate (in K per terrestrial day) produced by the CO<sub>2</sub> 15 $\mu$ m bands for the night time kinetic temperature profile.

## Heating rates

The solar heating rates show one major difference from Mars, in that the absorption coefficients decrease with height in the upper atmosphere because the Venus thermosphere is depleted in CO<sub>2</sub> above 140km. Between 90 and 110km bands of very different strengths give similar absorption rates, although peaked at different heights. The penetration of the solar flux into the atmosphere is larger at shorter wavelengths, where most of the weak bands are located, allowing them to make important contributions to the heating rate, as they do on Mars.

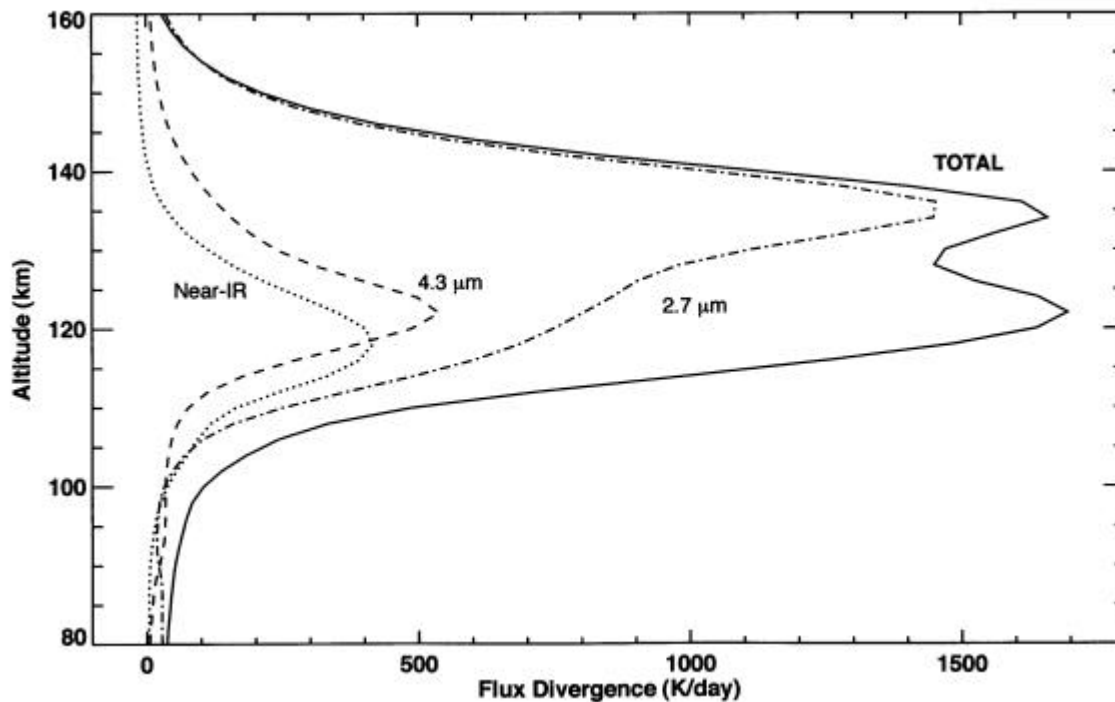


Figure 12. Total solar heating rates of the Venusian atmosphere for solar zenith angle  $=60^\circ$ . The total solar heating rate for a solar zenith angle of  $60^\circ$  appears in Fig.12, with the contributions from the different near-IR bands shown separately. As for Mars, much of the energy initially absorbed from solar radiation is emitted back to space, and radiative transfer and V-V collisions play major roles.

The maximum total solar heating amounts to more than 1500K/day between 120 and 135km with a double peak due to the different heights of the individual maxima, like the cooling rate curve. Significant cooling is found on Venus only above about 100km and below about 150km, the corresponding heights on Mars being 65 and 125km (Fig.10). Above these limits, all of the solar energy absorbed is re-emitted to space before it can be quenched by thermal collisions. In the lower atmosphere, the solar flux is too attenuated to have much effect.

### **Radiative equilibrium temperature**

In terms of global mean energy budget, carbon dioxide cooling at 15 microns dominates over eddy diffusion below 200 km. and over molecular conduction below 150 km; while solar heating in its near-IR bands is larger than the EUV heating below about 140 km. Therefore, from 140 km down to about 60 km (where vertical convection takes over as the principal vertical heat transfer mechanism) it is reasonable to apply a radiative equilibrium model to simulate the Venus mesosphere and lower thermosphere.

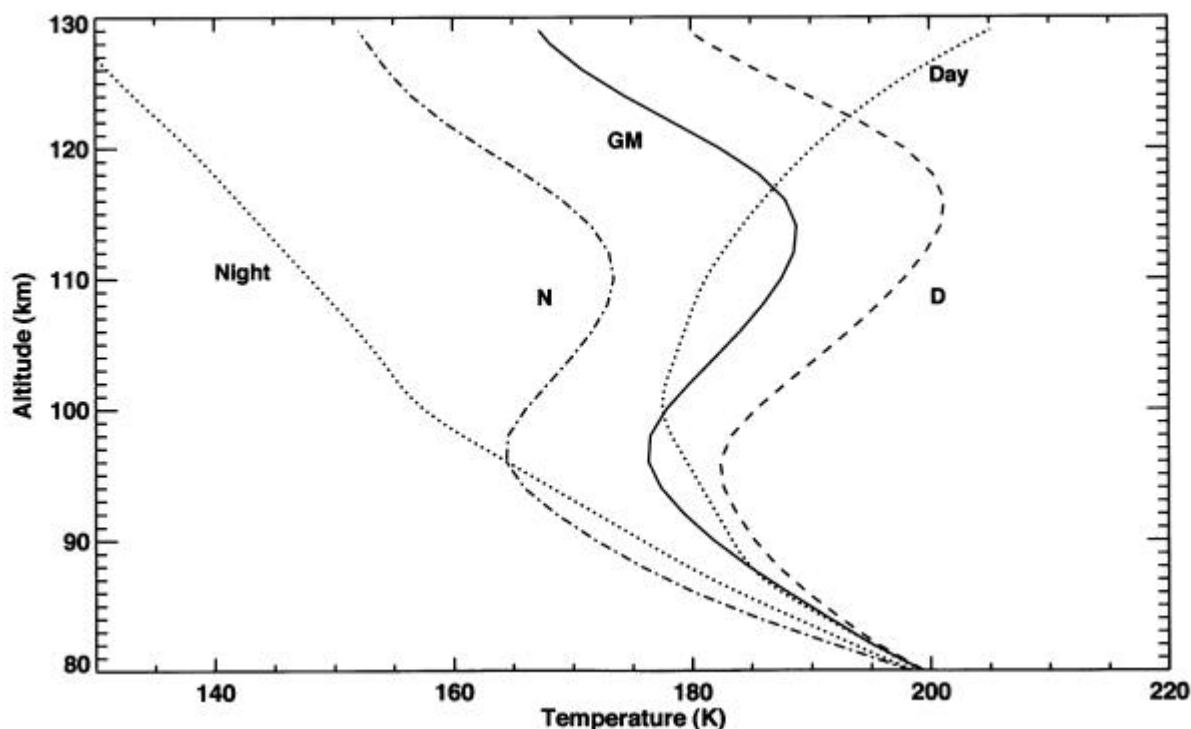


Figure 13. Radiative equilibrium temperature profiles for dayside (D), global mean (GM), and nightside (N). The reference temperature profiles are also plotted (dotted lines).}

The global mean radiative equilibrium temperature in Fig.13 was computed using a diurnally averaged solar heating rate, calculated for a solar zenith angle of 60deg and divided by 2. This shows a maximum around 110km produced by the solar absorption in the near-IR and in the weak bands at 4.3 and 2.7 $\mu$ m. The inflexion above 125km is due to the solar heating in the 2.7 $\mu$ m bands.

The differences between day and night-time conditions are particularly extreme, given the slow rotation of Venus. The purely radiative model predicts very cold temperatures in the night-time thermosphere, and a contrast with the daytime which is larger than Pioneer Venus observed. The difference is probably due primarily to adiabatic heating by dynamics, rather than errors in the radiative transfer calculations and their rate coefficients.

#### 5.4 Aeronomy

Aeronomy refers to the study of non-equilibrium processes which occur most prominently at low atmospheric densities. Frequently these are associated directly or indirectly with solar excitation and give rise to non-equilibrium radiative emissions that, under the generic name of airglow, can be detected and used to study the upper atmosphere. In addition to the intrinsic interest in themselves, these processes are sometimes (in the absence of in-situ data) the only access to some atmospheric regions. A number of reviews of what is known in this field have appeared recently. The aeronomy processes can be grouped under three main headings, as follows.

#### **Processes related to the physical separation and escape of atmospheric constituents.**

Solar Lyman-alpha radiation excite hydrogen, and the resonant atmosphere airglow has been used to obtain information about atomic hydrogen above 200 km. The values found require a large upward flux of H in the lower thermosphere at the sub-solar point, which seems to support an interhemispheric circulation. Apparently the escape flux of H decreases with solar activity, as they do the sputtering of oxygen atoms and pick up of ions by the solar wind. The

exosphere, where the scale height and the mean free path are about equal, is located at 210 and 154 km during day and night, respectively.

UV nightglow in gamma and delta bands of NO, excited by radiative recombination of N and O atoms, has been particularly valuable for constraining the circulation in the region between 115 and 150 km. The 2 hour offset in the night peak of NO emission toward the morning terminator indicates retrograde super-rotating zonal (RSZ) winds around 50 m s<sup>-1</sup> at those altitudes.

#### **Non-LTE radiative transfer processes.**

Similarly excited by solar radiation, this group includes mostly those constituent emissions in the IR part of the spectrum, and the most important are the CO<sub>2</sub> bands between 1.2 and 4.3 microns, CO(1-0) and (2-1) bands at 4.7 microns, and the O<sub>2</sub> infrared atmospheric band at 1.27 microns. Images taken from Earth of these O<sub>2</sub> bands in the night hemisphere show rapid changes in intensity that may be associated with variations in downwelling of atomic oxygen transported by the global circulation from the sunlit hemisphere. Also, observations of rotational lines of the CO bands (1-0) and (2-1) have been used to derive sub-solar to anti-solar (SS-AS) winds at the altitudes 125-145 and 100-110 km, respectively. These emissions result from direct solar fluorescence at 4.7 microns and from UV photolysis of CO<sub>2</sub>.

Apparently the SS-AS winds increase with altitude, and are much more stable over time than the RSZ winds, which show a very large variability in short and long term scales.

#### **Dissociation, ionisation, and photochemistry.**

A current problem is related to carbon dioxide recombination in the mesosphere. As on Mars, the atmosphere is poor in the photodissociation products O and CO, but on Mars we know the photochemistry of H<sub>2</sub>O participates in a cycle to reconstitute CO<sub>2</sub>. On Venus, however, the mesosphere appears to be very dry and it is necessary to look for alternative catalytic cycles.

Important links between photochemistry and radiation can be highlighted. The Venusian thermosphere is colder than the Martian case because of its larger atomic oxygen abundance, and it is very stable against variations in the solar flux because when this increases (solar rotation, solar cycle, etc) the atomic oxygen increases, and consequently the IR cooling by CO<sub>2</sub>.

## **6 THE GLOBAL CIRCULATION AND DYNAMICS**

As a result of the slow rotation of the planet, its near circular orbit and small obliquity, the underlying circulation of Venus' atmosphere is quite simple. The Sun is always above the equator, to within a couple of degrees, so the air warmed at low latitudes rises and moves towards the poles, where it cools and descends before returning equatorwards at lower altitudes.

QuickTime™ and a  
QuickDraw decompressor  
are needed to see this picture.

Figure 14. (a) Profiles of the zonal (east-to -west) and meridional (equator-to-pole) wind on Venus as measured by tracking the Pioneer Venus and the Venera 8,9,10, and 12 entry probes. Note the suggestion of global layering, consisting perhaps of stacked 'Hadley' circulation cells of considerable latitudinal extent.

On top of this simple picture are superimposed various complications. One that has been observed for many decades, but remains difficult to explain, is the global 'super-rotation', which manifests itself in cloud structure which moves rapidly around the planet in a direction parallel to the equator. The cloud markings, which appear with high contrast through an ultraviolet filter, have their origin at heights 50 or 60 km above the surface (where the pressure is of the order of 100 mb) and travel around the equator in 4 to 5 days, corresponding to speeds near  $100 \text{ m s}^{-1}$ . This is more than 50 times faster than the rotation rate of the surface below. Measurements of the winds below the clouds (**Fig.14**), and calculations (from temperature data) of the winds above the cloud tops, show that the zonal wind speed declines at higher and lower levels, reaching values near zero at about 100 km and near the surface respectively.

Direct measurements of the winds 1 m or so above the surface by the Russian landers Venera 9 and 10 found velocities of  $\approx 1 \text{ m s}^{-1}$ . Tracking of the Pioneer and Venera landers during their descent showed that there is a steady decrease with height from the  $100 \text{ m s}^{-1}$  or so observed in the ultraviolet markings near the cloud tops. Earth-based observers had earlier shown, by the measurement of Doppler shifted emission lines from atmospheric gases, that the cloud-tracked winds do, in fact, apply to mass motions, rather than the phase speed of waves as had also been suggested.

Attempts have been made to explain these high zonal wind speeds on Venus by several mechanisms, all of which fall into one of three main categories, viz. (i) the gravitational interaction of the Sun with the atmospheric tides, (ii) the overhead motion of the Sun in the sky (the 'moving flame' mechanism), and (iii) the upward transport of momentum from the surface. Currently prevailing opinion favours a version of mechanism (iii), in which



momentum from the solid planet is transported by waves whose interaction with the main flow is complex and in which the mean meridional circulation plays an important role. Parameterizations have been found which are able to produce large zonal velocities in dynamical models of the Venusian atmosphere, although of course this is not the same as saying that we understand the forcing or dissipation mechanisms responsible for the transfer of momentum from the surface to the cloud tops.

The cloud motions which trace the zonal winds also reveal some aspects of the pattern of the meridional circulation on Venus. As expected on the simple theoretical grounds outlined above, Hadley cells exist in each hemisphere. These are global-scale circulation cells characterised by rising motion all around one constant latitude belt and descending motion at another. Each cell extends to higher latitudes than on Earth, in part a consequence of the slower zonal rotation speeds. Near the poles on Venus, a complex instability develops, resulting in dramatic long-lived wave structures. The *polar collar* takes the form of a ribbon of very cold air, some 10 km deep and a thousand km in radius, centred on the pole. Inside the collar, temperatures are some 40 K cooler than outside the feature (**Fig. 15**). Poleward of the inner edge of the collar lies the *polar dipole*, a wavenumber 2 feature consisting of two well-defined warm regions circulating around the pole. Both the dipole and the collar have so far resisted attempts to model them as normal modes of the atmosphere.

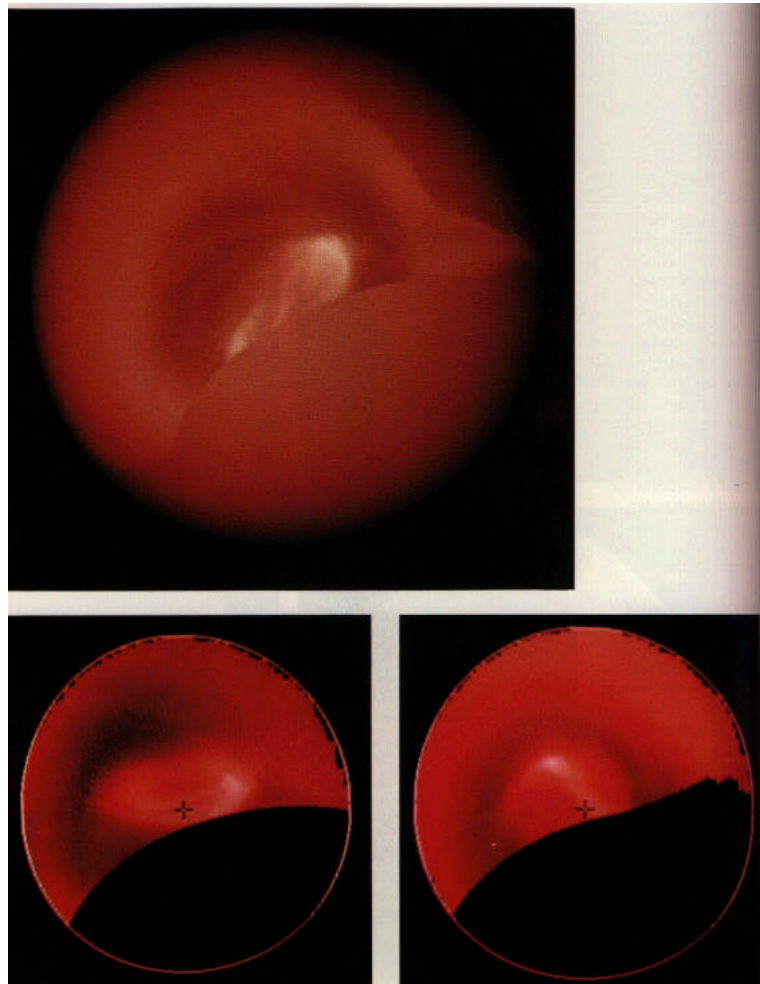


Figure 15. The polar dipole, imaged at 11.5  $\mu\text{m}$  in the thermal infrared by Pioneer Venus.

The cloud-tracked winds obtained from Pioneer Venus and Mariner 10 both show equator-to-pole velocities of around five metres per second in each hemisphere. Tracking of the PV probes shows winds of this magnitude at about 50 - 60 km altitude, with a very complicated

vertical structure (fig. 14). One possible interpretation of the alternations in the direction, as well as the magnitude, of the meridional wind, is that these could mark the passage of the probe through the different components of a stack of Hadley cells, each extending from the equator to high latitudes. The layered eddy sources and sinks which have been postulated as driving the zonal super-rotation may be related to the cell interfaces in this scenario.

Motions in the deeper atmosphere were observed by near-infrared imaging carried out by the Galileo probe in 1990. The features observed on the night side of the planet at wavelengths from 1 to 3.5  $\mu\text{m}$  originate in the main cloud deck, illuminated from below by thermal emission from the hot lower atmosphere. The typical velocities inferred near the equator were about half as fast as those from uv markings, which is consistent with the vertical profiles of wind and cloud opacity measured by the Pioneer and Venera probes, since the cloud layer providing most of the opacity in the case of the near-ir markings is 10 to 15 km deeper than for the UV markings. The Galileo winds feature a zonal jet of more than 100  $\text{ms}^{-1}$  at middle latitudes, and equator-to-pole drifts of a few  $\text{ms}^{-1}$ .

## **7 MODEL ATMOSPHERES**

### *7.1 A VENUS REFERENCE ATMOSPHERIC MODEL*

Most of the observational information we have about the Venus atmosphere comes from the Pioneer Venus (PV) mission of the late 1970s. More recently, high resolution infrared spectra were obtained by the Near-IR Mapping Spectrometer (NIMS) during the Galileo fly-by in 1991, but temperatures were derived only between 60 and 100km. Due to the significant day-night variation, we need to consider separate day and night-time reference temperature profiles, both taken from the PV measurements and shown in Fig.13. The densities and atomic oxygen abundances appear in Fig.16. Note that  $\text{CO}_2$  is the major constituent only below about 140km.

### *7.2 GLOBAL CIRCULATION MODELS*

#### **Basic Concept**

Global or general circulation models are a widely used tool in weather forecasting, climate forecasting and meteorological research. The basic idea is to produce a simulation of an atmosphere by solving the partial differential equations (PDEs) that govern the motion of the atmosphere subject to initial and boundary conditions. Although in principle it is possible to solve the equations exactly, there is no computer in the world that is powerful enough to do this. Also the initial conditions i.e. the state of an atmosphere at any particular time and the boundary conditions e.g. the surface albedo, cannot be known exactly. Hence one must approximate the system of PDEs and initial and boundary conditions in order to make the problem tractable. This approximation gives limitations on the performance of a GCM. Classical approaches to this problem and associated limitations, which afflict all GCMs, are described below.

#### **Dynamical formulations and discretization**

The motion of an adiabatic, frictionless planetary atmosphere can be described using the Navier-Stokes equations of fluid motion. When modelling an atmosphere however, it is appropriate to solve an approximate form of these equations namely the primitive equations (see e.g. Holton [1979]). Such an approximation is necessary, for instance, as sound waves which have high phase speeds which would require the time step over which the model is iterated to be very small. The primitive equations exclude sound waves so the model time step can be lengthened.

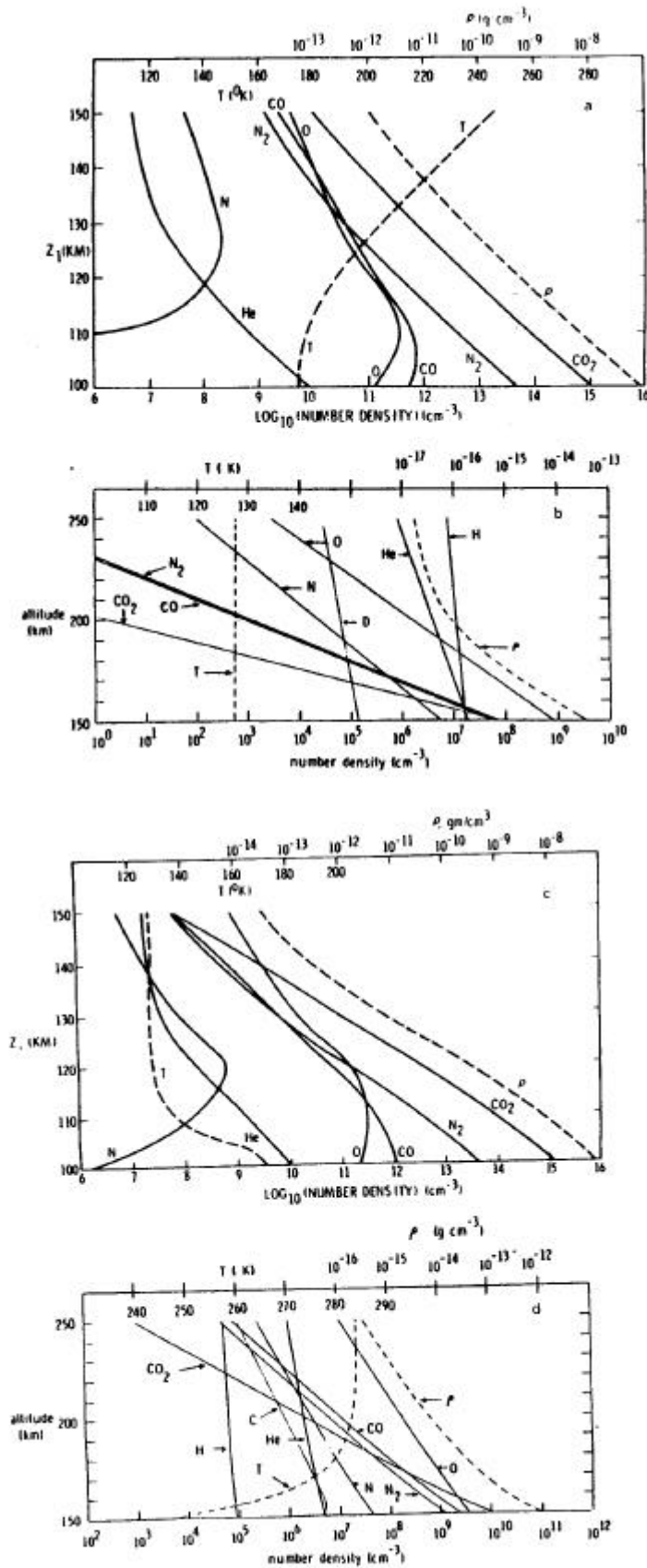


Figure 16. Volume mixing ratios of the main compounds of the Venusan atmosphere taken from the Venus International Reference Atmosphere (Keating et al., 1985) The top two frames are for daytime conditions with night-time below.

Once one has the primitive equations it is necessary to discretize them onto a grid of points in the horizontal and vertical. This discretization may be done any number of ways (see e.g. Haltiner and Williams [1979]) but inevitably it introduces errors into the solution of the PDEs. Firstly the equations of motion have many conservation properties e.g. the conservation of angular momentum, of which only a few can be guaranteed by a given numerical formulation. In fact it is interesting to note that both the GCMs proposed for use in this study do have numerical schemes that conserve angular momentum and this has been shown by Hourdin [1989] to be very important when simulating the Martian atmosphere. This property should be still more important for simulating Venusian atmospheric super-rotation.

Secondly, the limitations of computers require that solutions must be truncated in space and time, or in a spectral space for spectral models.

### **Parameterizations**

Because the equations of motion are represented on a finite number of grid points there are a number of processes e.g. turbulent eddies generated near the surface, that are not explicitly resolved. However these processes can be important in providing an accurate simulation as there may be a significant transfer of heat or momentum from so-called sub-grid scale processes, into the resolved scales of the GCM. Also effects such as the heating and cooling of the atmosphere due to radiative processes must be accounted for. Representing such things in a GCM falls under the general heading of parameterization. Parameterization schemes are often based on empirical formulas or idealised models and simulations of the terrestrial atmosphere are often very sensitive to the 'tuneable parameters' involved in such schemes.

### **Application to planetary atmospheres**

Although developed originally for terrestrial meteorology, physics-based GCMs can be applied in principle to other planetary atmospheres, at least for the terrestrial planets (solid bodies surrounded by a thin atmosphere). The first GCM of the Martian atmosphere was developed in the late sixties, during the very early history of terrestrial GCMs.

This pioneering work by Leovy and Mintz (1969) has been since followed by a number of other models developed around the world, which have confirmed the relevance of such tools for the Martian atmosphere.

When adapting a terrestrial GCM to Mars, the formulation of the dynamical equations and the associated numerical formulations are essentially unchanged. The main effort needed is on the "parameterizations", in particular on radiative transfer. The conditions in which radiation is absorbed and emitted by the thin CO<sub>2</sub> and dusty atmosphere of Mars can not be handled with exactly the same radiative codes as those developed for our wet and cloudy atmosphere. Most recent Martian GCMs such as the LMD-Oxford Model also include parameterizations of physical processes specific to the atmosphere of the red planet: CO<sub>2</sub> condensation and sublimation, radiative impact of CO<sub>2</sub> snow, dust-lifting and transport, and so on. GCMs now play a central role in the study of the Martian atmosphere and climate as they do on Earth.

Past attempts to develop specific GCMs for Venus have been quite limited. The only full treatment published to date is that by Young and Pollack (1977). The main result of their model was the simulation of strong prograde equatorial winds compatible with the observed super-rotation of Venus. However, this super-rotation was obtained by including in the model an ad hoc parameterization accounting for equatorward transport of angular momentum by eddies. Similar efforts conducted at LMD in the early 80s concluded that super-rotation can be maintained and even reinforced without such a parameterization, but these remained unpublished (Tourte, 1984).

Although no other GCM has been developed specifically for Venus, the ability of GCMs in general to simulate atmospheric super-rotation has been shown for idealized planetary

conditions (Hourdin et al., 1991, Del Genio et al., 1993) and for Titan (Hourdin et al. 1995). This question, crucial for the development of a new GCM for Venus, is detailed in the next section. We then discuss what would be the main work to be done on existing GCMs such as the LMD-Oxford Mars Model to develop a reliable model of the atmosphere of Venus, suitable for compiling a dedicated database.

### *7.3 SIMULATING ATMOSPHERIC SUPER-ROTATION WITH GCMs*

The dynamics of terrestrial atmospheres is driven primarily by latitudinally varying insolation. A classical view suggests that such a forcing should result in strong prograde winds in the middle or high latitudes as the consequence of poleward angular momentum transport by a thermally direct Hadley circulation. Those winds could be called super-rotating in the sense that, locally, the air rotates in the same direction but faster than the solid planet. However, the super-rotation of Venus (and probably Titan) is far from this picture. The atmosphere of Venus, at 70 km above the surface, rotates globally 60 times faster than the planet. For Titan, the mid stratosphere, near 250 km, probably rotates ten times faster than the surface.

In both cases, even the air above the equator is strongly super-rotating. Such a super-rotation can not be explained by a simple axi-symmetric redistribution of angular momentum. In particular, equatorial super-rotation, requires the presence of non-axisymmetric eddy motions (Hide 1969, Held and Hou 1980, Rossow 1985). A number of mechanisms for producing such eddies have been proposed, including barotropic instability of the high-latitude jet produced by the Hadley cell (Gierasch 1975, Rossow and Williams 1979), solar semidiurnal thermal tides (cf. Fels and Lindzen 1974, Pechmann and Ingersoll 1984, Leovy 1987), transient or topographically forced planetary- or small-scale gravity waves (Leovy 1973, Hou and Farrell 1987, Gierasch 1987, Del Genio and Rossow 1990), and external torques (Gold and Soter 1971).

GCMs predict the existence of high-latitude jets when run at slow rotation, but were for many years unsuccessful in generating equatorial super-rotation. As stated above, the GCM of Young and Pollack (1977) did produce strong equatorial winds, but in the presence of unrealistic momentum sources, due to a non-conservative eddy diffusion parameterization.

First attempts to produce super-rotation by running GCMs for slowly rotating planets have yielded weak or even subrotating zonal winds at the equator despite the presence of quasi-barotropic eddies which transport momentum equatorward as required by the Gierasch-Rossow-Williams scenario (Rossow 1983, Covey et al. 1986, Del Genio and Suozzo 1987). Since none of the latter group of models included a diurnal cycle, it was plausible to speculate that the missing transport by thermal tides is the crucial super-rotation mechanism. Williams (1988a,b) in fact achieves moderate equatorial super-rotation via diurnal forcing of a moist terrestrial GCM. This is appropriate for the upper Venus atmosphere, but much stronger than that experienced by the deep atmosphere. Consistent with this view, Hou et al. (1990), using a high resolution tidal model coupled to a zonally averaged circulation model, have demonstrated that tides act primarily within and above the cloud layer on Venus and can account for only about half of the observed super-rotation.

#### **Recent simulations of atmospheric super-rotation**

The first attempt to simulate super-rotation concentrated on the slow rotation of the solid planet. In fact, another condition seems to be quite important: the stability of the vertical temperature profile. Del Genio and Suozzo (1987) suggested that the failure of slowly rotating GCMs to super-rotate was a consequence of excessive vertical mixing by parameterized convection driven by statically unstable radiative equilibrium temperature profiles.

Hourdin et al. (1992) first succeeded in producing super-rotation without any ad hoc forcing.

They used a "parametric" version of the LMD GCM in which the description of the planet was reduced to a set of 20 parameters. Whereas this model was able to produce a reasonable Earth-like meteorology for terrestrial values of those parameters, a strong super-rotation was obtained with different values. In addition to the slow rotation rate, a key in this simulation was the use of a strong coefficient for the absorption of solar radiation which significantly enhanced the stability of the atmosphere.

Del Genio et al (1993) obtained strong super-rotating regimes using a terrestrial general circulation model (GCM), simplified by removal of the hydrologic cycle, the diurnal cycle, and all seasonal and geographic variations, to study a suite of equilibrium simulations in which rotation rate, radiative heating profile, surface drag, and stratospheric drag are varied. The key to super-rotation in these experiments is the presence of an upper troposphere cloud which intercepts much of the incoming solar flux and produces a statically stable radiative equilibrium state in the lower and middle troposphere. The radiative heating profile limits the depth of boundary layer convection and detaches the upper level flow from the surface. The cloud-covered GCM produces equatorial winds in excess of 50/100 m sec<sup>-1</sup> with/without stratospheric drag. In these simulations, removal of the cloud decreases static stability, increases vertical convective mixing, and almost completely eliminates equatorial super-rotation.

A significant super-rotation was also obtained more recently in a Titan GCM (Hourdin et al., 1995). Those simulations are also characterized by the slow rotation rate of the planet and by the strong absorption of solar radiation by the orange haze within the atmosphere.

So, an underestimation of vertical stability, resulting in an overestimation of turbulent mixing or convection seems to be the main reason for the failure in simulating super-rotation. Indeed, the convection mixes momentum downward where it is transferred to the solid planet by surface drag, and the adiabatic thermal structure prevents large vertical wind shears. The resulting weakly stable lapse rate inhibits the development of quasi-barotropic eddies. In slowly rotating terrestrial GCMs, the radiative state is dictated by the semi-transparency of Earth's atmosphere to sunlight, which is absorbed mostly at the surface. In the Venus GCM of Rossow (1983), the radiative equilibrium state was specified to be super-adiabatic in the lowest 20 km, according to prevailing opinion about the Venus thermal structure at that time. Pioneer Venus, Venera, and Vega probes have since shown, however, that the Venus atmosphere is statically stable above 5 km altitude except for isolated near-neutral layers at ~25-30 km and ~50-55 km (cf. Young et al. 1987).

In all of the simulations mentioned here, super-rotation extends throughout most of the atmosphere (as it does on Venus). Hadley cell interaction with quasi-barotropic eddies which transport momentum equatorward is responsible for the excess angular momentum, in line with certain aspects of proposals by Gierasch (1975) and Rossow and Williams (1979) albeit with relatively modest Prandtl numbers. However, Limaye (1988a,b, 1990) has suggested, on the basis of analyses of the movement of UV cloud markings, that the observed eddy momentum transport, at the cloudtop level at least, may be in the opposite direction to that implied by the models.

In all these simulations, topography was not included, and in all but the Titan simulation, the diurnal and seasonal cycles and all geographic variations were removed. Each of these may actually play an important role on Venus. What those simulations do show, is that without them, a strong super-rotation can be reached just by the Gierasch-Rossow-Williams mechanism. Nevertheless, the mechanisms responsible for the creation of super-rotation in the model have still to be confirmed by observations of Venus and of Titan.

Another concern is that the GCM results could depend partly on the numerical approach employed. Del Genio et al. (1993) point out that an important feature of the GISS GCM is that

its stable numerical schemes permit the model to be run with no explicit horizontal or vertical diffusion. The Arakawa differencing techniques conserves global angular momentum to within about 0.0005% per day, or 1-10% over an entire integration of the model to equilibrium, with the numerics acting as a very weak momentum sink. This assures that any super-rotation which occurs is the result of explicitly resolved and/or parameterized physical processes.

On the other hand, studies performed with the LMD GCM (Hourdin et al., 1992, 1995) did include a parameterized super-rotation, and showed that this parameterization is crucial. The numerical scheme for large scale dynamics is built so as to ensure the conservation of potential enstrophy (enstrophy is the square of absolute vorticity). This property controls the interaction between different scales in 2D or quasi 2D turbulence, and results in a cascade of enstrophy from the large scale where it is injected, to the small scales where it is dissipated by viscosity. Since this cascade is properly represented in the GCM, the cascade systematically produces an accumulation of enstrophy at the smallest resolved scale of the model. Dissipating this enstrophy at the mesh scale is both physically realistic and essential for the model stability. The LMD simulations showed that this parameterization was not dominant for equatorward angular momentum transport, which was essentially produced by explicitly resolved barotropic planetary waves.

All together, those results suggest that GCMs are able to produce atmospheric super-rotation, following the Gierasch-Rossow-Williams mechanism, and that they will generally do so for slowly rotating planets with a stable enough atmosphere.

#### 7.4 DEVELOPING A NEW GLOBAL CIRCULATION MODELS FOR VENUS

This section describes possible steps in developing a new GCM specifically for Venus. Since simulations and models of the atmosphere of Venus are not yet available, this has to be quite speculative, and we suggest an incremental approach, with a progressive improvement of the GCM.

##### **General requirements**

Since the model is intended to produce useful information for future space missions, including landing and aero-assistance, it should probably cover the atmosphere from the surface up to about 140 km, as for the LMD-Oxford Mars GCM.

The horizontal discretization will probably be imposed by the need to simulate atmospheric super-rotation by explicitly resolving the important eddies, without being too strongly affected by parameterized dissipation. Our experience suggests that a spatial resolution of 3 by 3 degrees, similar to that used currently for Mars, should be enough.

The question of vertical resolution is probably more crucial since the entry profiles of the Venera probes suggest a layering of winds with a typical vertical length scale of order one km. It is difficult to postulate *a priori* if these small scale vertical structures must be accounted for in the GCM to obtain a realistic climate. Our experience however suggests that, if the grid is not fine enough to capture those modes, it could result in the excitation of spurious numerical modes at the grid-scale. This suggests that the model should use a rather fine vertical resolution of the order of 1 km or a number of layers of the order of one hundred.

The model could start with a rather simple set of parameterizations for subgrid-scale motions, at least for turbulent mixing and dry convection. Such schemes are already included in the LMD-Oxford Mars GCM and can theoretically be reused without major changes.

The main outstanding question is that of how to deal with the computation of radiative transfer; this will be central in the incremental approach we propose.

## Radiative transfer

As explained earlier in this document, computing radiative transfer in the atmosphere of Venus is quite challenging. Conditions encountered vary from strong absorption in a very dense CO<sub>2</sub> atmosphere to quasi conservative scattering by clouds at visible and near infrared wave-lengths. As on Mars, the upper atmosphere requires special treatment to allow for non-LTE effects.

Even line-by-line codes, generally accurate schemes which have been developed for the purpose of interpreting observations of the Earth's atmosphere, must be used carefully under the novel conditions of high temperature and pressure on Venus. The latter have effects on the shape of individual spectral lines which are not well characterized at the present time. Perhaps the greatest problem lies in the derivation of a fast, numerically tractable code for a GCM in which radiative transfer can be computed every few minutes if necessary, over a broad spectral range and for vertical profiles associated with each point of the horizontal grid. The rest of the discussion assumes that a line-by-line reference model exists for the atmosphere of Venus and concentrates on the problem of deriving a code fast enough for a GCM.

The question of solar heating could be treated in a quite simple way as a first step. Indeed, if variations of obliquity, and atmospheric composition are neglected, the solar heating in the atmosphere only depends upon latitude and altitude. The thermal infrared is more difficult to handle since, because of the fundamental dependency of emission on atmospheric temperature, the computation can not be made a priori.

The net exchange formulation, developed for the LMD-Oxford Martian GCM may be a good way of approaching this problem. The idea of this method is to reformulate radiative transfer in terms of net exchange of energy between pairs of model layers rather than in terms of radiative fluxes. It can be shown quite easily that, for small enough spectral intervals, and for two isothermal layers, the net power exchanged between layer *i* and *j* can be expressed formally as

$$P_{i,j} = X_{i,j} (B_i - B_j)$$

where *B* is the average Planck function on the spectral interval and *X<sub>i,j</sub>* is a coefficient which only depends upon the optical properties of the atmosphere. *X<sub>i,j</sub> B<sub>i</sub>* is that part of the power emitted by layer *i* which is absorbed by layer *j*. If both layers are at the same temperature, this power is exactly canceled by that emitted by *j* and absorbed by *i*.

For two nearby layers, assuming that layers are isothermal introduces errors. However, a linear temperature profile can also be handled by the simple formula above. Because it ensures the conservation and reciprocity properties of radiative transfer are not violated it can be used safely to introduce approximations into the computations. In the case where there is a combination of absorption and scattering, as on Venus, the determination of coefficient *X<sub>i,j</sub>* as a function of optical path and spectral properties can be very complex.

This net exchange formulation could lead to the following three steps in developing a radiative code for Venus.

*Step 1:* If latitudinal and temporal variations of optical properties are neglected, the exchange coefficients *X<sub>i,j</sub>* could be estimated numerically from a series of computations performed with a line-by-line code on a reference temperature profile. For each computation, the source function would be set to zero in all layers except layer *i*. With such a source function, the line-by-line computation would directly produce the net power *P<sub>i,j</sub>* and hence optical coefficients *X<sub>i,j</sub>*.

This approach requires *N* times *M* integrations of the line-by-line code where *N* is the number of layers retained for the GCM and *M* is the number of spectral intervals used (typically of the order of a few hundreds).



Although very demanding, this computation can be done once for all. Exchange coefficients are then stored and reused during the GCM simulation in each point of the horizontal grid. The method is straightforward and has already been tested for Venus a few years ago at LMD. Rerunning and assessing this work will represent an important part of the development of a new Venus GCM.

*Step 2:* In step 1, the vertical pressure grid must be the same in the line-by-line code and in the GCM. In particular, it is directly applicable only to a billiard-ball Venus. The inclusion of orography, which results in a varying pressure discretization depending on the location on the surface of the planet, would require that we are able to interpolate the exchange coefficient from one vertical grid to another. This step has not been adequately tested yet. The solar heating can also be computed once and for all. Provided it is done as a function of longitude, latitude and altitude.

*Step 3:* The last step would be to be able to account for possible variations of the cloud cover or atmospheric composition. This requires the means to be able to compute on-line, in the GCM, the dependency of the exchange coefficients upon the atmospheric composition. There are probably two main directions to develop such a model.

The first one consists of extracting the required information from additional integrations of the line-by-line model. For instance, the sensitivity of the exchange coefficient to the cloud opacity could be stored as well and used in the GCM simulation. The second way would consist in developing a fast, approximate radiative transfer code, based for instance on k-distribution methods, which could be used regularly in the GCM to update the exchange coefficient, with a periodicity depending upon the time-scale of the physical processes which modify the atmospheric composition and optical properties. The choice between approaches will require a detailed analysis of these processes, which include cloud micro-physics and possibly chemistry. These may have to be included in the model as well. This step would also require some work on the part of the code that deals with visible wavelengths.

This description in three steps mainly concerns the part of the atmosphere where LTE is valid. The non-LTE part will have to be treated in a different way. However, the treatment of non-LTE for the Martian atmosphere has much in common with that required for Venus. The recent comprehensive study by Roldan et al. (2000) will be useful to guide our work and that we anticipate that an adaptation from the Martian parameterizations to the Venus case should not introduce any special difficulties.

## **8 A VENUS CLIMATE DATABASE**

The EuroMars consortium has gained considerable experience during the development of the Mars Climate Database (Lewis et al. 1999) in summarizing GCM output and making it available to the wider community. Novel techniques have been derived to represent the variability which is inevitably lost in a statistical representation of GCM output and to add variability on scales smaller than the GCM is able to represent explicitly, e.g. internal gravity waves. Once a realistic Venus GCM has been developed, we plan to apply the same methodology to produce a Venus Climate Database based directly on its output.

The Venus database will follow the structure of the MCD, although it may be unnecessary to provide the same sort of interannual variability (dust scenarios) that was necessary for Mars. This will have to be assessed once experience is gained with a Venus GCM. The small obliquity of Venus also means that the division of the database into 12 seasons, as was done for Mars, may also be unnecessary. Both of these factors will allow the resulting database to be fitted into a much smaller amount of disk space.

On the other hand, local solar time of day is clearly an important factor in the Venusian atmosphere and it may well prove desirable to store fields at more local times of day (12 were

used for Mars; this may be at least doubled for Venus). The thick atmosphere is also likely to mean that the GCMs will be run with more vertical levels and this should be reflected in the database structure. The net effect of these changes may be to produce a database of comparable size, with the mean and variance of atmospheric fields stored as functions of longitude, latitude, height and local time of day. Different levels of solar activity may also have to be considered for the upper atmosphere.

The large-scale (empirical orthogonal function) and small-scale (gravity wave) variability models are equally applicable to Venus, with the observational evidence (summarized earlier) indicating the presence of both small- and planetary-scale waves. The 2D EOF model, as used in the MCD versions 2 and above, can be retained to represent the large-scale variability in the Venus GCM, with EOFs calculated from frequent model output to capture as much of the variance as possible. Small-scale gravity wave motions are also likely to be of great importance, and the scheme used for Mars will have to be re-calibrated for Venus.

## **9 CONCLUSIONS**

This document summarises most of current knowledge concerning the surface and atmospheric environment of Venus. In the next stage of the project we will produce a numerical model of the Venusian atmosphere which draws on this information to produce a realistic simulation of the atmosphere, and we will use this model to produce a database of statistics of the atmosphere. Please see the model and database user and software requirement documents for more information.

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